

Past, Present and the Future of Precision $\sin^2 q_W$ Measurements

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*NuFact'03, June 9, 2003
Columbia University, New York*

- Introduction
- Past measurements
- Current Improvement
- Measurement at a Neutrino Factory
- Conclusions

NuTeV Collaboration

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$\sin^2 q_W$ and n-N scattering

- In the electroweak sector of the Standard Model, it is not known *a priori* what the mixture of electrically neutral electromagnetic and weak mediator is → This fractional mixture is given by the mixing angle
- Within the on-shell renormalization scheme, $\sin^2 \theta_W$ is:

$$\sin^2 q_w^{\text{On-Shell}} = 1 - \frac{M_W^2}{M_Z^2}$$

- Provides independent measurement of M_W & information to pin down M_{Higgs}
- Comparable size of uncertainty to direct measurements
- Measures light quark couplings → Sensitive to other types (anomalous) of couplings
- In other words, sensitive to physics beyond SM → New vector bosons, compositeness, ν -oscillations, etc

How did we measure?



$$\text{coupling} \propto I_{\text{weak}}^{(3)}$$

$$\text{coupling} \propto I_{\text{weak}}^{(3)} - Q_{EM} \sin^2 \theta_W$$

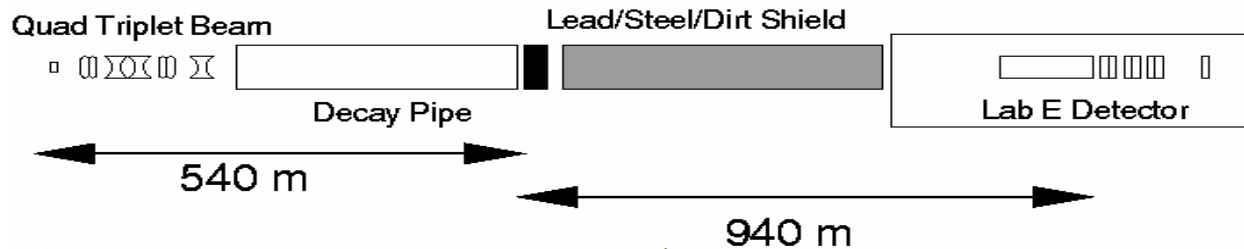
- Cross section ratios between NC and CC proportional to $\sin^2 \theta_W$
- Llewellyn Smith Formula:

$$R^{n(\bar{n})} = \frac{S_{\text{NC}}^{n(\bar{n})}}{S_{\text{CC}}^{n(\bar{n})}} = ?^2 \left(\frac{1}{2} - \sin^2 \theta_W + \frac{5}{9} \sin^4 \theta_W \left(1 + \frac{S_{\text{CC}}^{\bar{n}(n)}}{S_{\text{CC}}^{n(\bar{n})}} \right) \right)$$

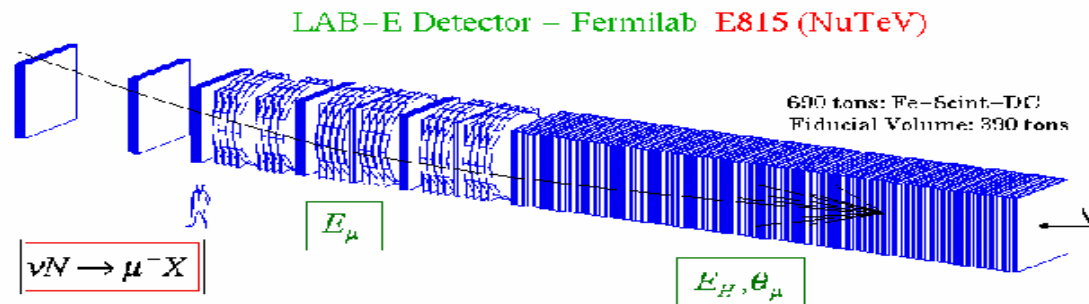
Some corrections are needed to extract $\sin^2 \theta_W$ from measured ratios (radiative corrections, heavy quark effects, isovector target corrections, HT, R_L)

Previous Experiment

E770: Quad Triplet Beam and Lab E Detector



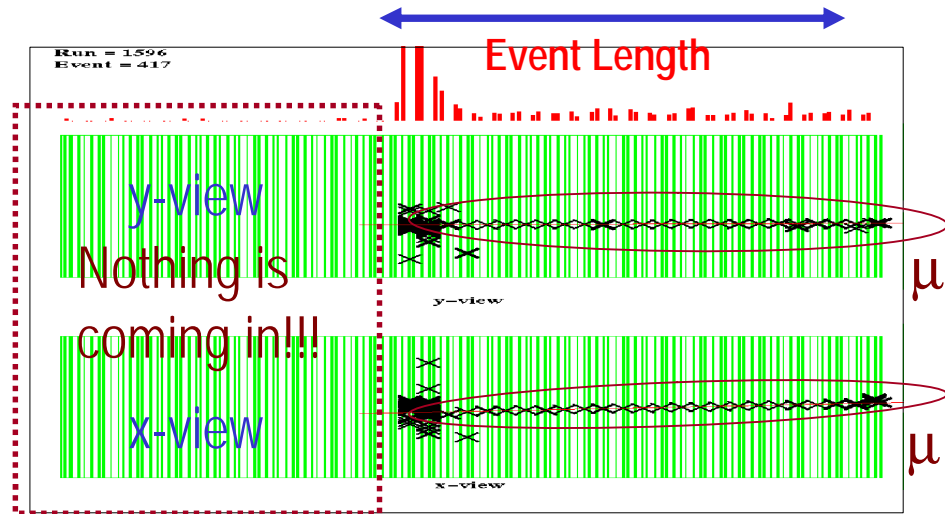
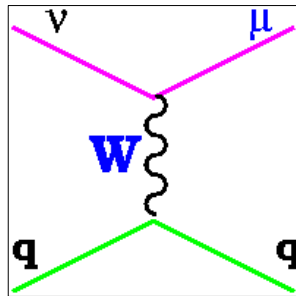
- Conventional neutrino beam from π/k decays
- Focus all signs of π/k for neutrinos and antineutrinos
- Both $\bar{\nu}_\mu, \nu_\mu$ in the beam (NC events are mixed)



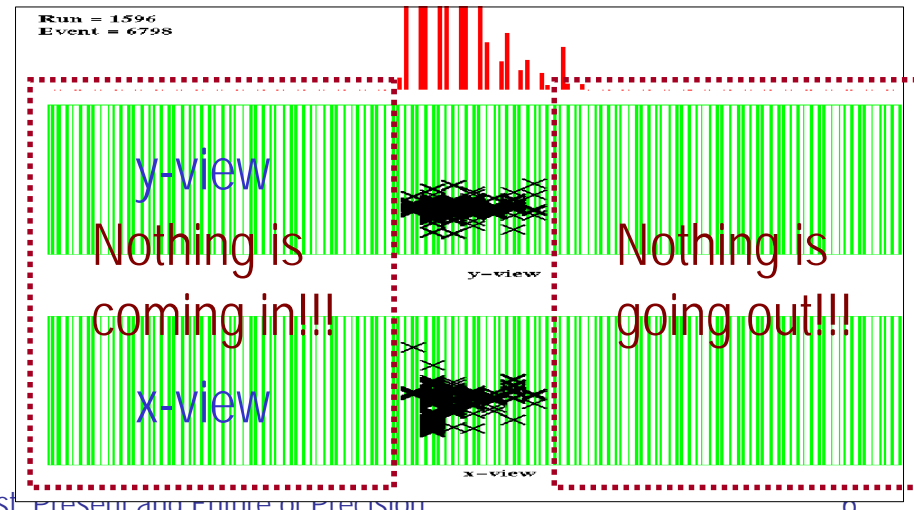
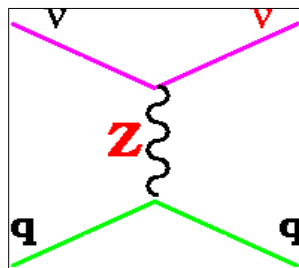
- Very small cross section \rightarrow Heavy neutrino target
- ν_e are the killers (CC events look the same as NC events)

How Do We Separate Events?

Charged Current Events



Neutral Current Events



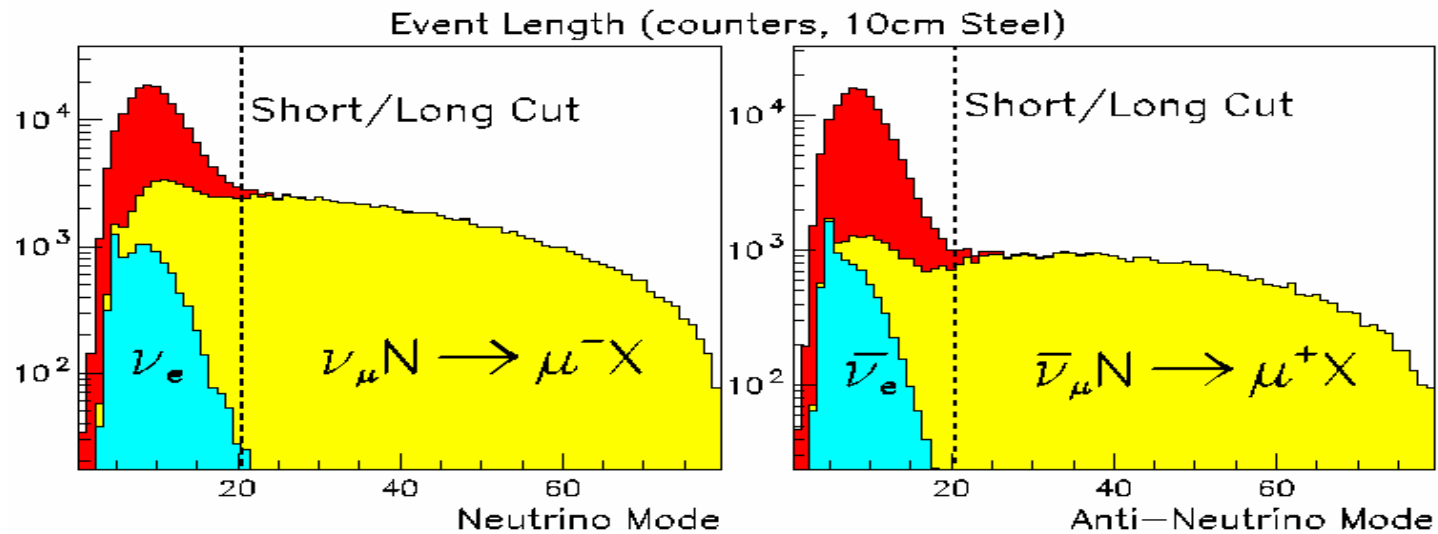
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Event Length

Define an Experimental Length variable

→ Distinguishes CC from NC experimentally in statistical manner



Compare experimentally measured ratio

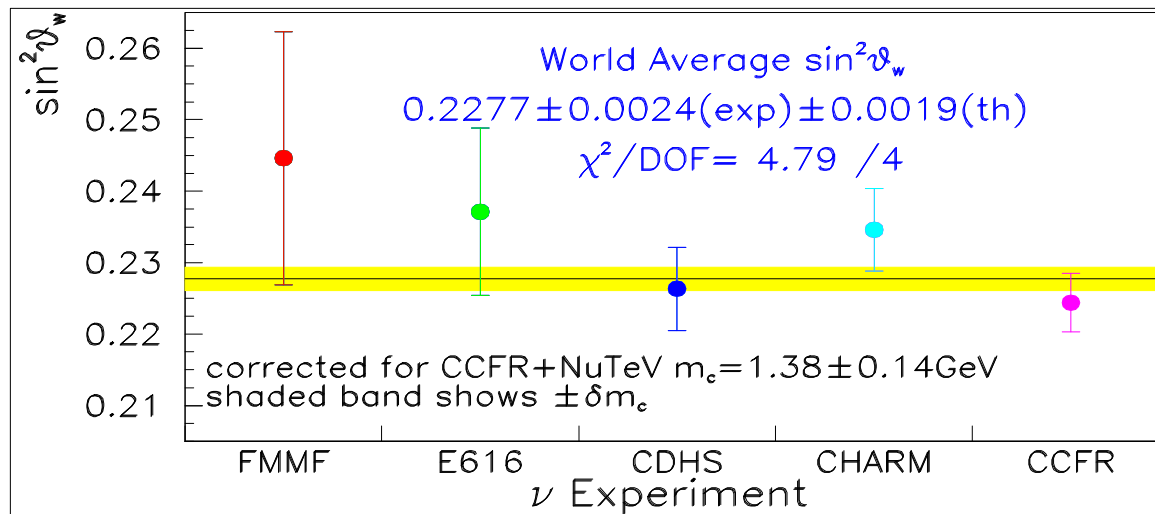
$$R_{\text{Exp}} = \frac{N_{\text{Short}}}{N_{\text{Long}}} = \frac{L < L_{\text{Cut}}}{L > L_{\text{Cut}}} = \frac{N_{\text{NC Candidates}}}{N_{\text{CC Candidates}}}$$

to theoretical prediction of R^ν

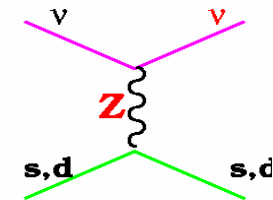
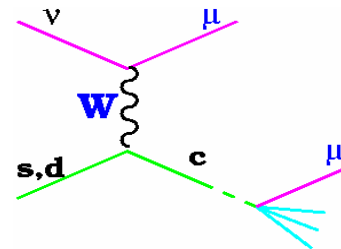
Past Experimental Results

$$\sin^2 \theta_W^{\text{On-Shell}} = 1 - \frac{M_W}{M_Z} = 0.2277 \pm 0.0031$$

$$\Rightarrow M_W^{\text{On-Shell}} = 80.14 \pm 0.16 \text{ GeV}/c^2$$



- Significant correlated error from CC production of charm quark (m_c) modeled by slow rescaling, in addition to ν_e error



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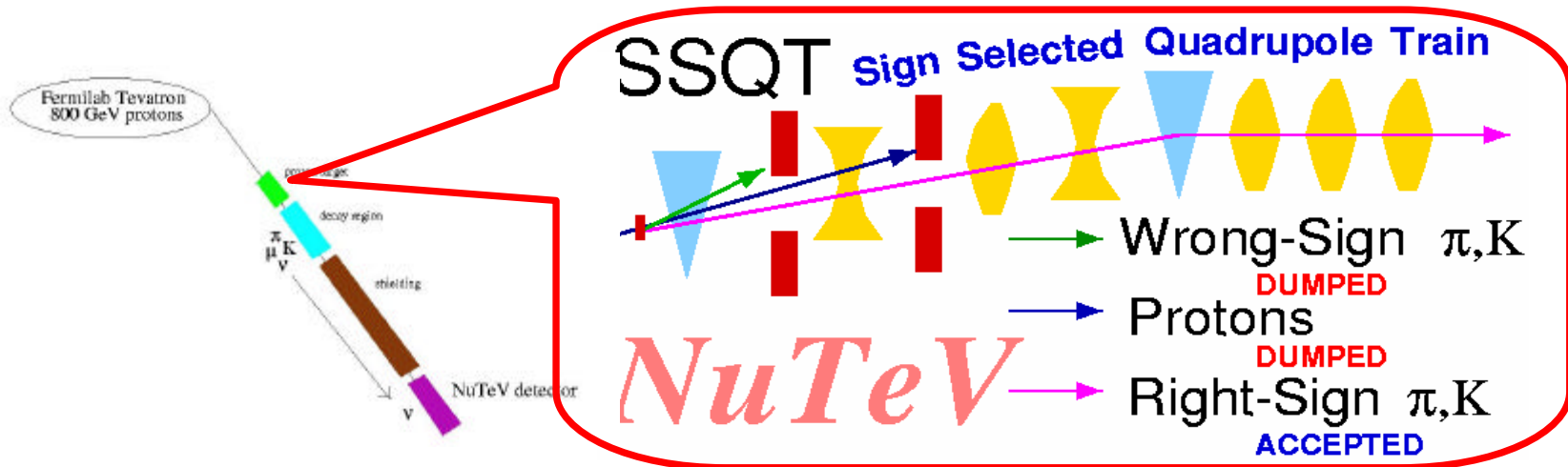
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The NuTeV Experiment

- Suggestion by Paschos-Wolfenstein by separating ν and $\bar{\nu}$ beams:

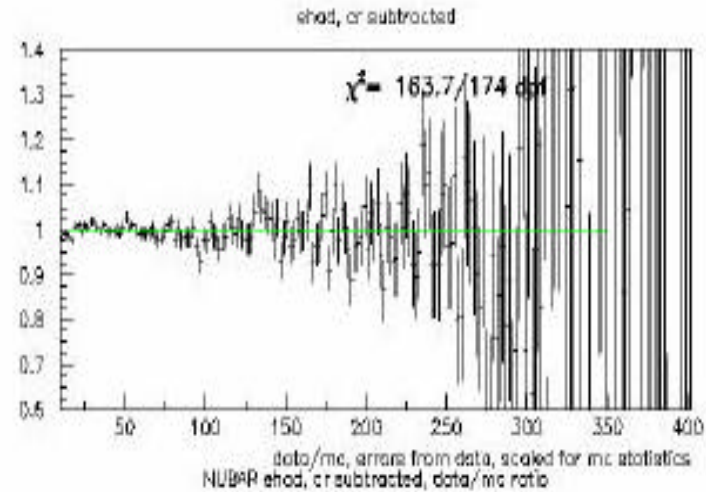
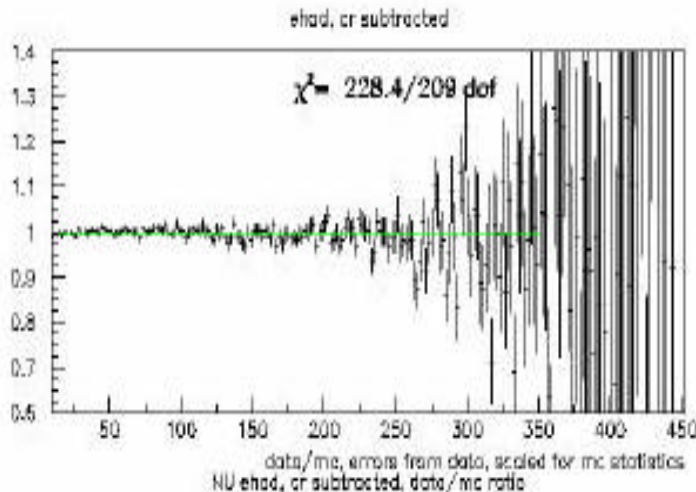
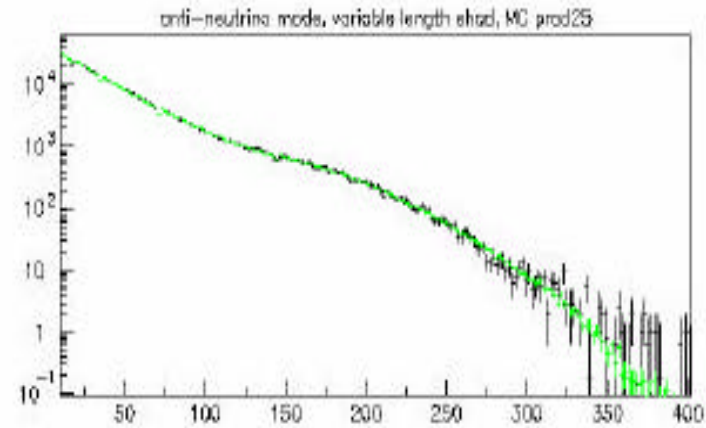
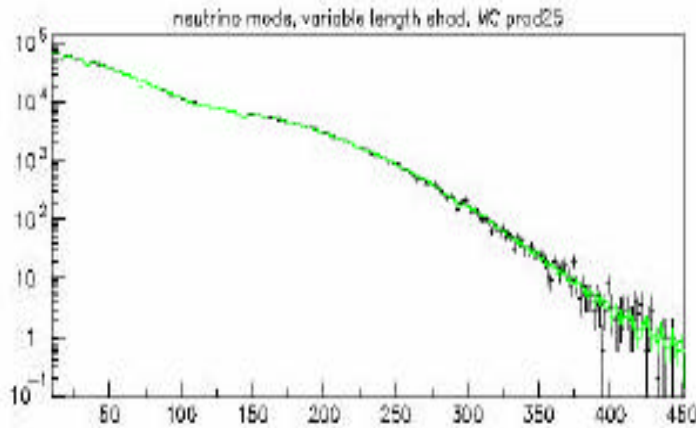
$$R^- = \frac{S_{NC}^n - S_{NC}^{\bar{n}}}{S_{CC}^n - S_{CC}^{\bar{n}}} = ?^2 \left(\frac{1}{2} - \sin^2 \theta_w \right) = \frac{R^n - R^{\bar{n}}}{1-r}$$

- Reduce charm CC production error by subtracting sea quark contributions
 - Only valence $u, d,$ and s contribute while sea quark contributions cancel out
 - Massive quark production through Cabibbo suppressed d_s quarks only
- Smarter beamline
 - Separate ν and $\bar{\nu}$ beam
 - Removes all neutral secondaries to eliminate ν_e content



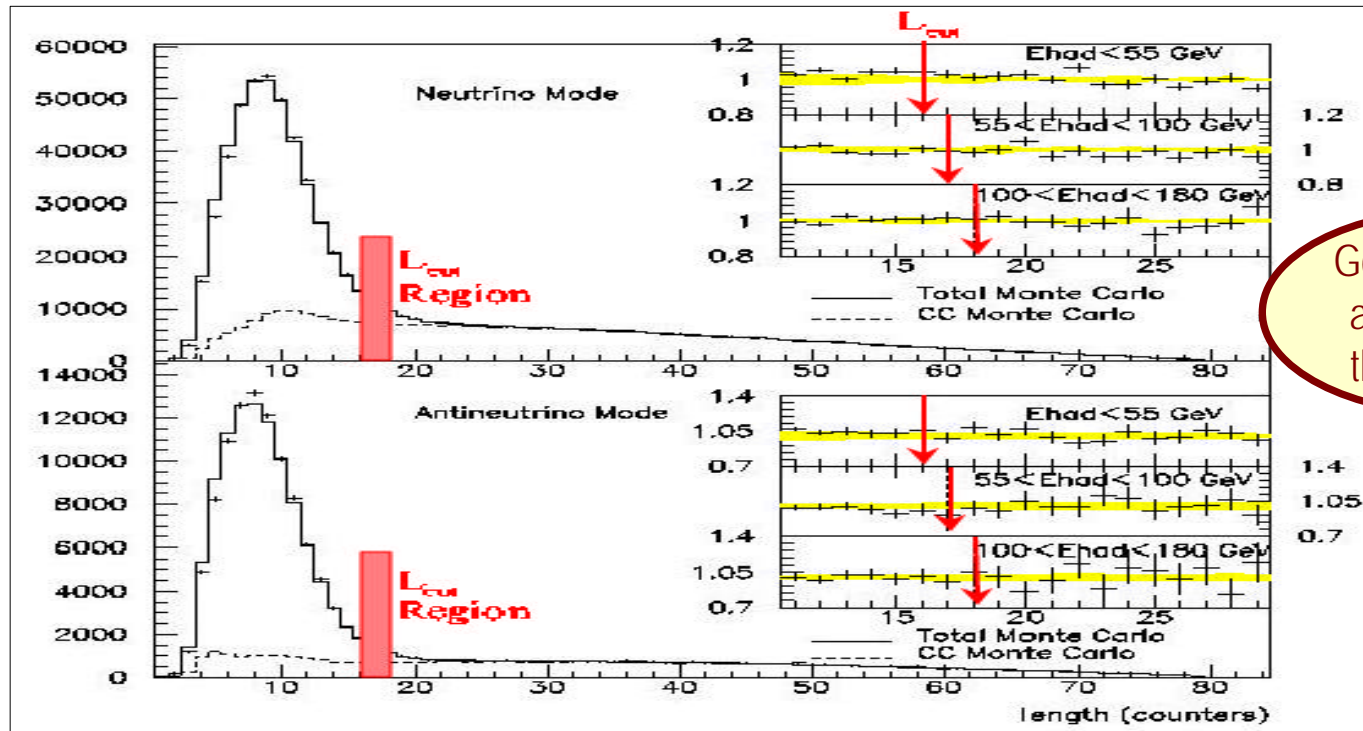
Events and E_{Had} After Event Selection

Events passing cuts: $1.62\text{M } \nu$ & $350\text{k } \bar{\nu}$ ($\langle E_{\nu} \rangle \sim 100\text{GeV}$)



NuTeV Event Length Distributions

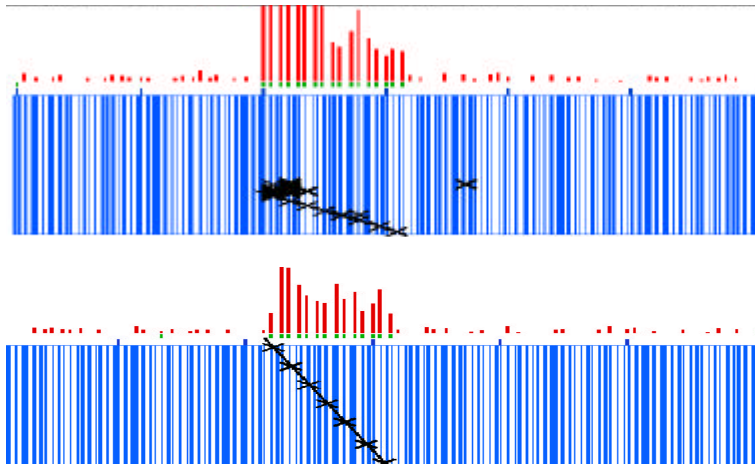
Energy Dependent Length cut implemented to improve statistics and reduce systematic uncertainties.



Good Data-MC agreement in the cut region

Mode	N_{short}	N_{long}	$R_v = N_{\text{short}}/N_{\text{Long}}$
n	457k	1167k	$0.3916 \pm 0.0007(\text{stat})$
\bar{n}	101k	250k	$0.4050 \pm 0.0016(\text{stat})$

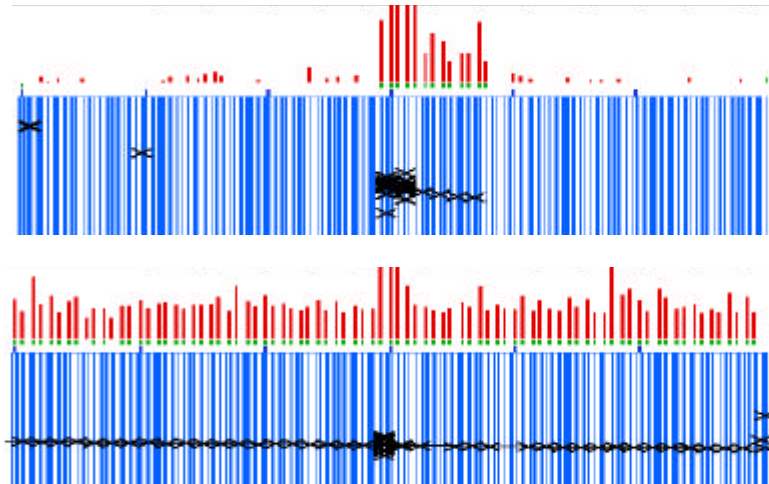
Event Contamination and Backgrounds



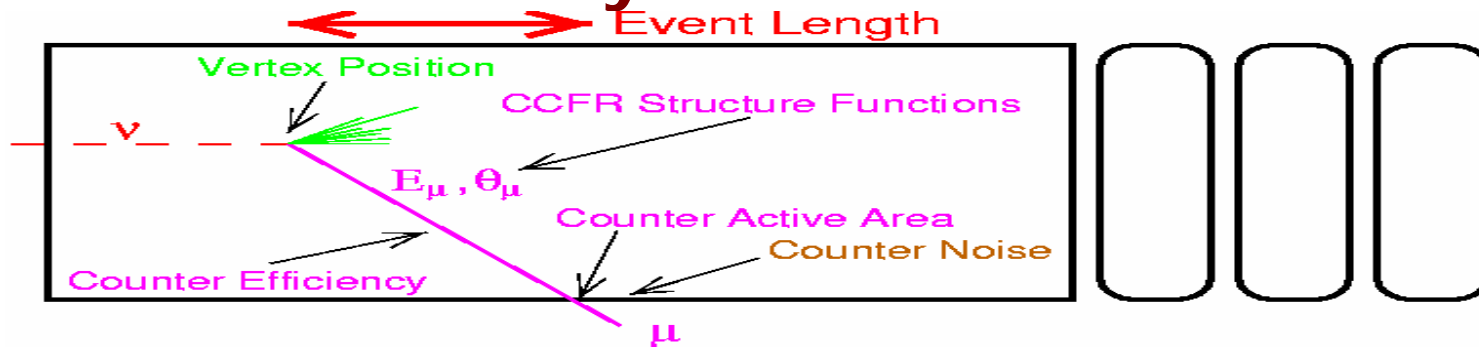
- SHORT n_m CC's (20% n , 10% \bar{n})
 μ exit and rangeout
- SHORT n_e CC's (5%)
 $\nu_e N \rightarrow eX$
- Cosmic Rays (0.9%)

- LONG n_m NC's (0.7%)
hadron shower
punch-through effects

- Hard m Brem(0.2%)
Deep μ events



Other Systematic Effects



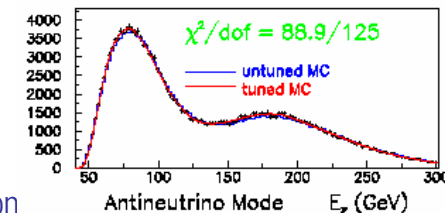
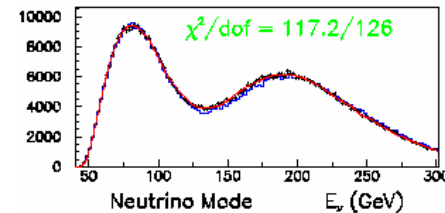
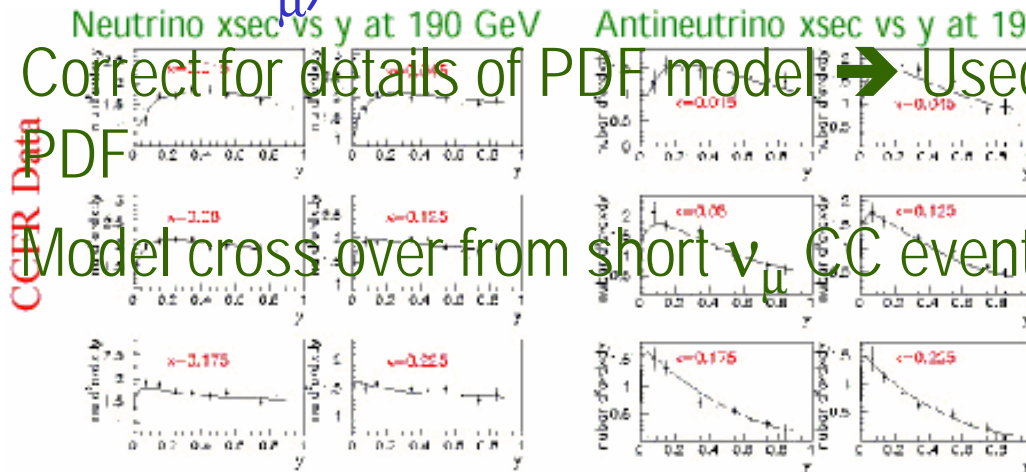
Sources of experimental uncertainties kept small, through modeling using n and TB data

Effect	Size($d\sin^2q_w$)	Tools
Z_{vert}	0.001/inch	$\mu^+\mu^-$ events
X_{vert} & Y_{vert}	0.001	MC
Counter Noise	0.00035	TB μ 's
Counter Efficiency	0.0002	ν events
Counter active area	0.0025/inch	ν CC, TB
Hadron shower length	0.0015/cntr	TB π 's and k 's
Energy scale	0.001/1%	TB
Muon Energy Deposit	0.004	ν CC

MC to Relate R_n^{exp} to R^n and $\sin^2\theta_W$

- Parton Distribution Model ($\langle Q^2 \rangle \sim 25 \text{ GeV}^2$ for ν_μ , 16 GeV^2 for $\bar{\nu}_\mu$)

- Correct for details of PDF model \rightarrow Used CCFR data for PDF
- Model cross over from short ν_μ CC events



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sin²θ_W Fit to R_n^{exp} and R _{\bar{n}} ^{exp}

- Thanks to the separate beam → Measure R^v's separately
- Use MC to simultaneously fit R_n^{exp} and R _{\bar{n}} ^{exp} to sin²θ_W and m_c, and sin²θ_W and ρ

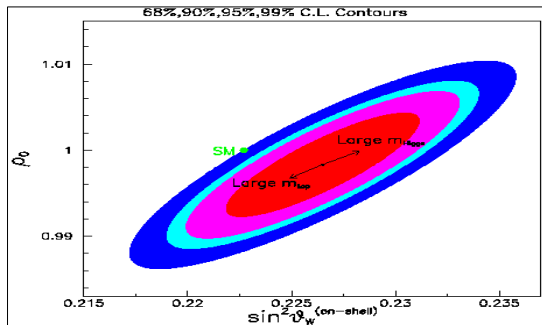
$$R^{n(\bar{n})} = \frac{S_{NC}^{n(\bar{n})}}{S_{CC}^{n(\bar{n})}} = ?^2 \left(\frac{1}{2} - \sin^2 \theta_W + \frac{5}{9} \sin^4 \theta_W \left(1 + \frac{S_{CC}^{\bar{n}(n)}}{S_{CC}^{n(\bar{n})}} \right) \right)$$

- R^v Sensitive to sin²θ_W while R ^{\bar{v}} isn't, so R^v is used to extract sin²θ_W and R ^{\bar{v}} to control systematics
- Single parameter fit, using SM values for EW parameters (ρ₀=1)

$$\sin^2 \theta_W = 0.2277 \pm 0.0013 \text{ (stat)} \pm 0.0009 \text{ (syst)}$$

$$m_c = 1.32 \pm 0.09 \text{ (stat)} \pm 0.06 \text{ (syst) w/ } m_c = 1.38 \pm 0.14 \text{ GeV}/c^2 \text{ as input}$$

- Two parameter fit for sin²θ_W and ρ₀ yields



$$\sin^2 \theta_W = 0.2265 \pm 0.0031$$

$$\rho_0 = 0.9979 \pm 0.041$$

Syst. Error dominated since we cannot take advantage of sea quark cancellation

NuTeV $\sin^2 q_W$ Uncertainties

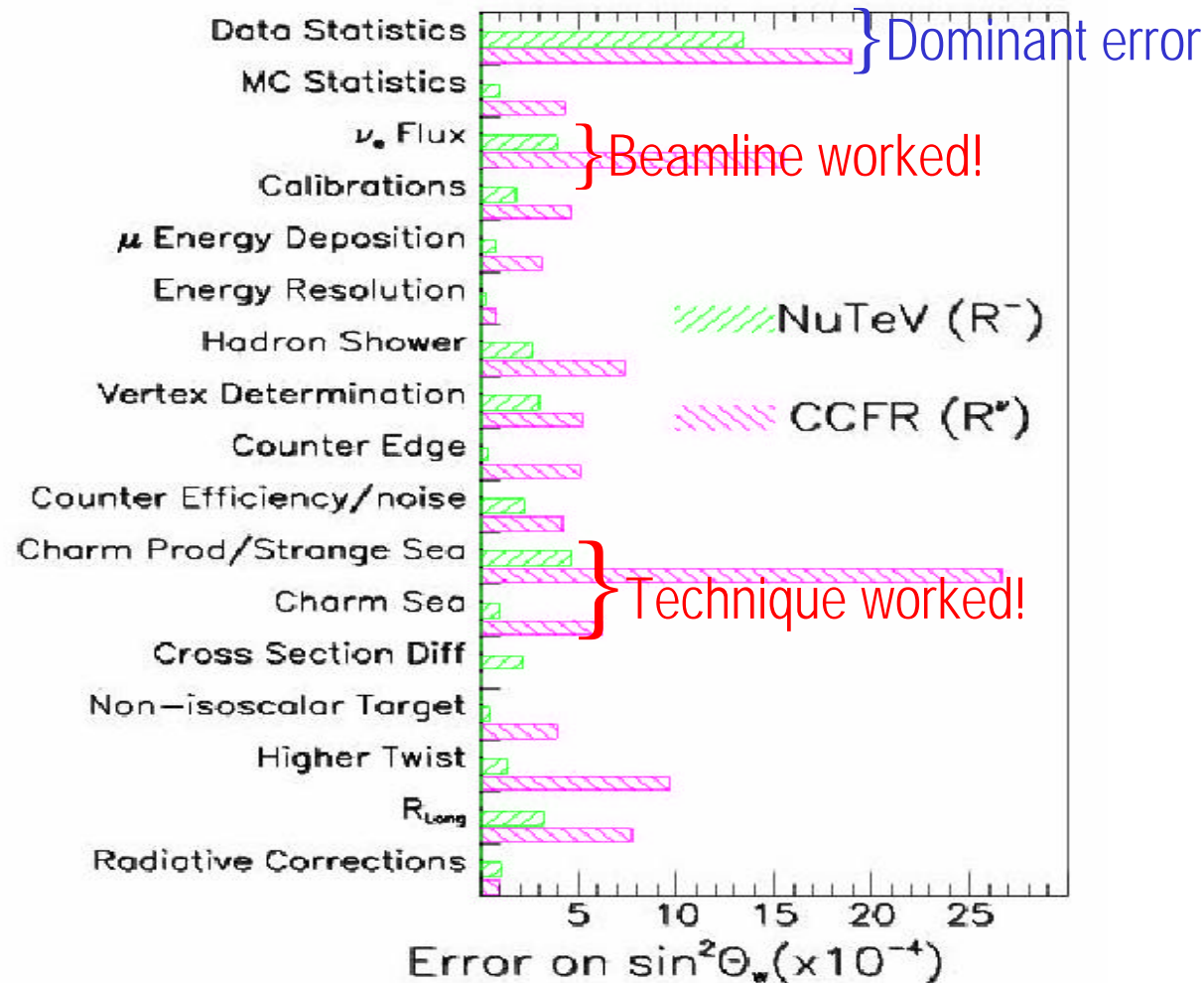
Source of Uncertainty	$d \sin^2 q_W$
Statistical	0.00135
ν_e flux	0.00039
Event Vertex	0.00030
Length (Other effects)	0.00027 (23)
Total Experimental Systematics	0.00063
CC Charm production, sea quarks	0.00047
R_L	0.00032
$s^{\bar{n}} / s^n$	0.00022
Higher Twist	0.00014
Radiative Correction	0.00011
Non-isoscalar target	0.00005
Total Physics Model Systematics	0.00064
Total Uncertainty	0.00162
DM_W (GeV/c²)	0.08

Dominant uncertainty

1-Loop Electroweak Radiative Corrections based on Bardin, Dokuchaeva **JINR-E2-86-2 60 (1986)**

$$d \sin^2 q_W^{(On-shell)} = -0.00022 \times \left(\frac{M_t^2 - (175 \text{ GeV})^2}{(50 \text{ GeV})^2} \right) + 0.00032 \times \ln \left(\frac{M_H}{150 \text{ GeV}} \right)$$

Past vs Present Uncertainty Comparisons



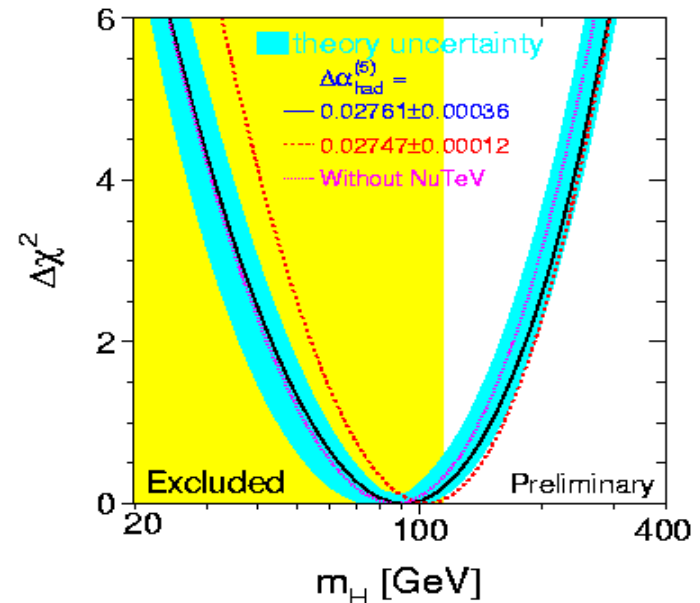
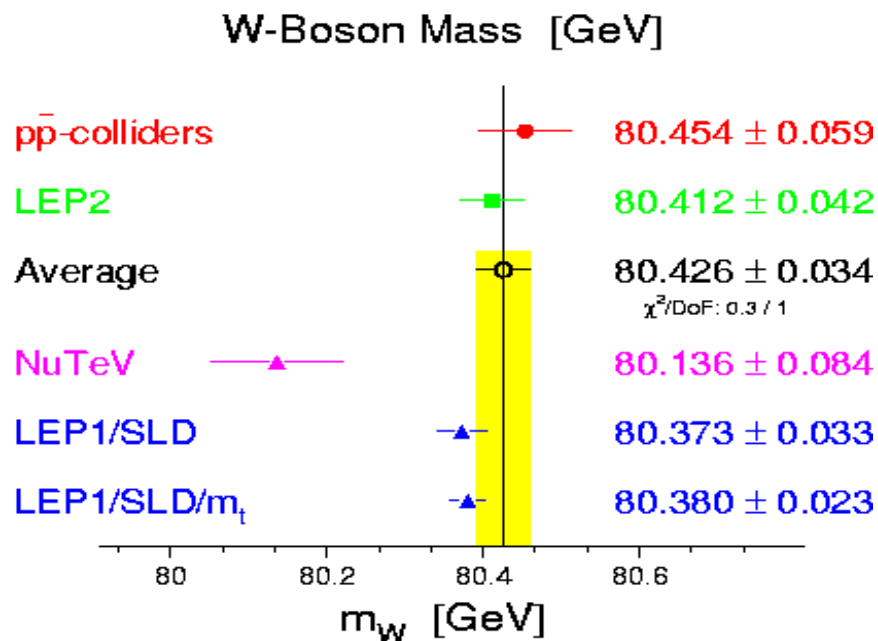
The Present (NuTeV) $\sin^2\theta_W$

$$\sin^2\theta_W^{\text{On-shell}} = 0.2277 \pm 0.0013 \text{ (stat)} \pm 0.0009 \text{ (syst)}$$

$$\sin^2\theta_W^{\text{On-shell}} = 1 - \frac{M_W^2}{M_Z^2}$$

$$\Rightarrow M_W^{\text{On-shell}} = 80.14 \pm 0.08 \text{ GeV/c}^2$$

Comparable precision but value smaller than other measurements



Confidence level in upper M_{higgs} limit weakens slightly.

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Model Independent Analysis

- Performed the fit to quark couplings (and g_L and g_R)

- For isoscalar target, the νN couplings are

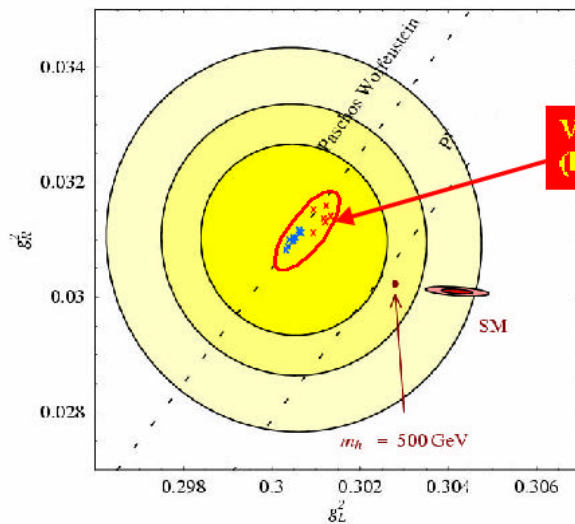
$$g_L^2 = u_L^2 + d_L^2 = g_0^2 \left(\frac{1}{2} + \sin^2 \theta_w + \frac{5}{9} \sin^4 \theta_w \right)$$

$$g_R^2 = u_R^2 + d_R^2 = g_0^2 \frac{5}{9} \sin^4 \theta_w$$

- From two parameter fit to R_n^{exp} and R_n^{exp}

$$g_L^2 = 0.3001 \pm 0.0014 \quad (\text{SM: } 0.3042 \leftarrow -2.6\sigma \text{ deviation})$$

$$g_R^2 = 0.0308 \pm 0.0011 \quad (\text{SM: } 0.0301 \leftarrow \text{Agreement})$$



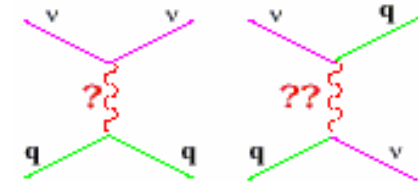
Difficult to explain the disagreement with SM by:
 Parton Distribution Function or LO vs NLO or
 Electroweak Radiative Correction: large M_{Higgs}

What is the discrepancy due to (Old Physics)?

- R- technique is sensitive to q vs \bar{q} differences and NLO effect
 - Difference in valence quark and anti-quark momentum fraction
- Isospin symmetry assumption might not be entirely correct
 - Expect violation about 1% → NuTeV reduces this effect by using the ratio of ν and $\bar{\nu}$ cross sections → Reducing dependence by a factor of 3
- s vs \bar{s} quark asymmetry
 - s and \bar{s} needs to be the same but the momentum could differ if +30% asymmetry
 - NuTeV LO di- μ measurement shows $\Delta s = s - \bar{s} \sim -0.0027$
 - NuTeV NLO analysis show no-asymmetry (D. Mason, et al., ICHEP02 proceedings)
- NLO and PDF effects
 - PDF, m_c , Higher Twist effect, etc, are small changes
- Heavy vs light target PDF effect (Kovalenko et al., hep-ph/0207158)
 - Using PDF from light target on Iron target could make up the difference → NuTeV result uses PDF extracted from CCFR (the same target)

What other explanations (New Physics)?

- Heavy non-SM vector boson exchange: Z' , LQ, etc
 - LL coupling enhanced than LR needed for NuTeV
- Propagator and coupling corrections
 - Small compared to the effect
- MSSM : Loop corrections wrong sign and small for the effect
- Gauge boson interactions
 - Allow generic couplings → Extra Z' bosons???
 - LEP and SLAC results says $< 10^{-3}$
- Many other attempts in progress but so far nothing seems to explain the NuTeV results
 - Lepto-quarks
 - Contact interactions with LL coupling (NuTeV wants $m_{Z'} \sim 1.2\text{TeV}$, CDF/D0: $m_{Z'} > 700\text{GeV}$)
 - Almost sequential Z' with opposite coupling to ν



Langacker *et al*, Rev. Mod. Phys. **64** 87; Cho *et al.*, Nucl. Phys. **B531**, 65;
 Zppenfeld and Cheung, hep-ph/9810277; Davidson et al., hep-ph/0112302

NLO Upgrade of $\sin^2\theta_W$ Analysis

- To address concerns within the community
 - Don't expect to see large effects
 - LO x-sec model describe CC x-sec data well
 - Gambino, et al., (hep-ph/0112302) shown little NLO PDF effect to R- σ level shifts to R- small (Davidson, hep-ph/0112302 & Kulagin, hep-ph/0301045)
- To calculate $O(\alpha_s)$ pQCD corrections to the differential X-sec for ν , $\bar{\nu}$ DIS
- NuTeV (Zeller & McFarland) is collaborating with Theorists
 - FNAL Theory group: K. Ellis, B. Dobrescu, W. Gigele
 - DESY: Seven-Olaf Moch
 - others
- Approach based on Altarellu, Ellis & Martinelli, NP B143, 521 (1978)
 - X-sec's written in terms of xF_1 , F_2 , xF_3
 - pQCD corrections affect $2xF_1-F_2=F_L$ & xF_3-F_2
 - F_L effect taken into account via R_L
 - Need α_s correction of xF_3-F_2 , because α_s^2 is small (Zijlstra, PLB297, 377, 1992)
- Calculations and Implementation of the correction in progress

$\sin^2 q_w$ Measurement at a NuFact

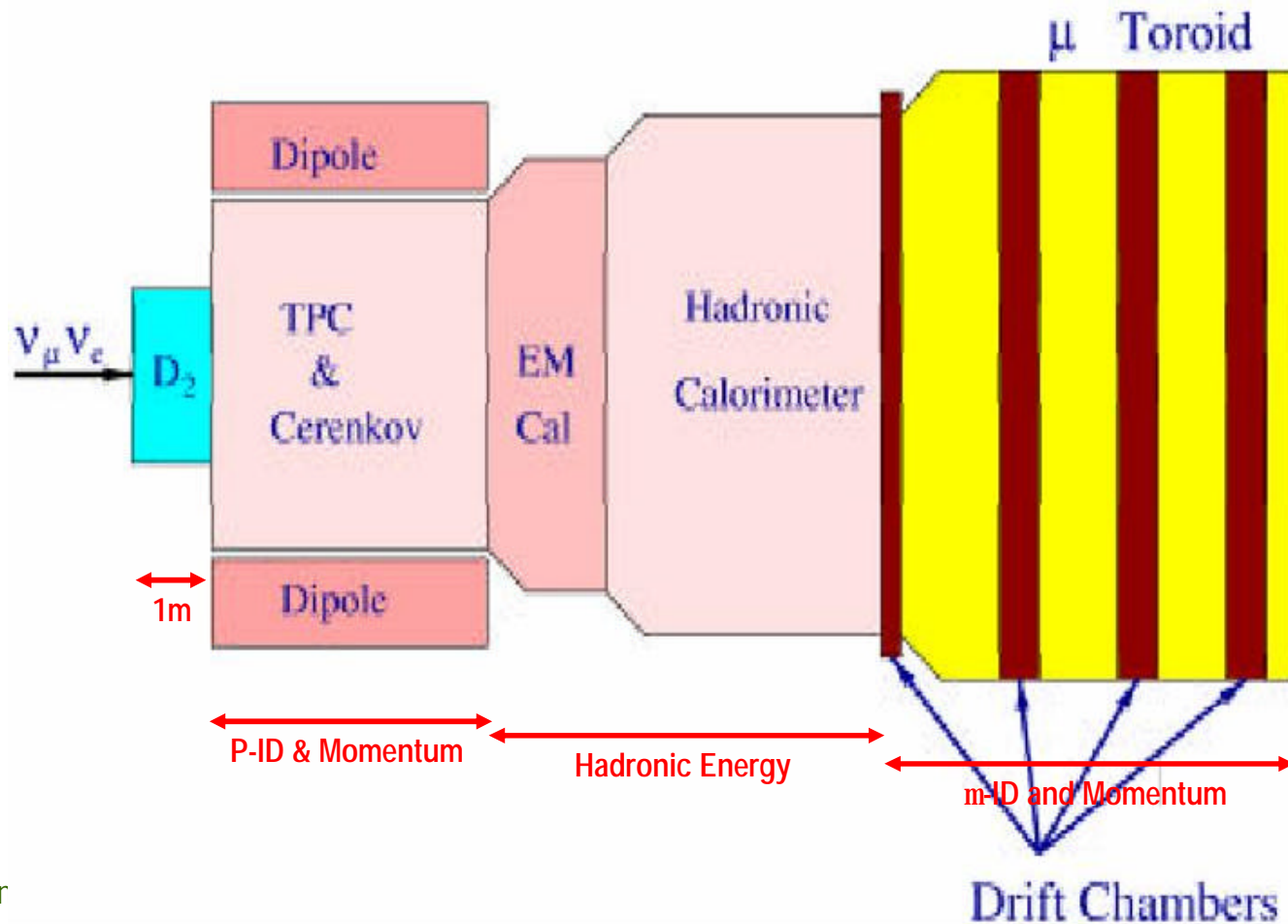
- Neutrinos come from μ decays
 - Good understanding of the beam content and flux
- Better collimated than conventional beam
- Large neutrino flux ($10^5 \sim 10^6$ higher than the current)

But...

- Always two neutrinos simultaneously in the given beam
($\nu_e + \nu_\mu$ or $\nu_\mu + \nu_e$)
 - Traditional heavy target detector will not work
 - Will screw up NC counting due to ν_e CC events
 - Need light target detectors → Can afford to do this
- Might need new techniques for NC to CC ratio
 - Can't distinguish ν_e vs ν_μ induced NC events

A Light Target $\sin^2\theta_W$ Detector at a NuFact

ν_e and ν_μ from muon decays are in the beam at all times
→ Must use light target (D_2) detectors



Expectation at a NuFact

Using a 1m thick D₂ target, one can obtain about 20M ν_μ CC events per year → With the help of good p-id, the stat doubles → Length related uncertainties become irrelevant

Source of Uncertainty	$d \sin^2 q_w$
Statistical	$1.35 \times 10^{-3} \rightarrow 2.13 \times 10^{-4}$
ν_e flux	$3.9 \times 10^{-4} \rightarrow 0$
Event Vertex	$3.0 \times 10^{-4} \rightarrow 3.0 \times 10^{-6}$
Energy Measurements	$1.80 \times 10^{-4} \rightarrow 9.00 \times 10^{-5}$
Total Experimental Systematics	$6.30 \times 10^{-4} \rightarrow 9.00 \times 10^{-5}$
CC Charm production, sea quarks	$4.70 \times 10^{-4} \rightarrow 2.40 \times 10^{-4}$
Higher Twist	$1.40 \times 10^{-4} \rightarrow 1.40 \times 10^{-4}$
Non-isoscalar target correction	$5.00 \times 10^{-5} \rightarrow 0$ (D ₂ target)
s^+ / s^-	$2.20 \times 10^{-4} \rightarrow 1.50 \times 10^{-4}$
Radiative Correction	$1.10 \times 10^{-4} \rightarrow 1.10 \times 10^{-4}$
R_L	$3.20 \times 10^{-4} \rightarrow 9.00 \times 10^{-5}$
Total Physics Model Systematics	$6.40 \times 10^{-4} \rightarrow 4.6 \times 10^{-4}$
Total Uncertainty	$1.62 \times 10^{-3} \rightarrow 5.15 \times 10^{-4}$
DM_W (GeV/c²)	$0.08 \rightarrow 0.025$

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Experimental and Theoretical Issues

Experimental Issues

- Must be able to reverse beam polarity and measure current well
- Detector must be light weight
- Must be able to distinguish primary e , μ , and π
 - Need to control overall p-ID efficiency to be better than 10^{-3}
- High electron detection efficiency
- Good EM and Hadronic shower ID
- Good charged particle momentum measurement
- Good vertex measurement w/ triggering capability at the target

Theoretical Issues

- Better measured charm CC x-sec
- Need to understand radiative correction better
- Better understanding of higher twist effects

Conclusions

- NuTeV has improved $\sin^2\theta_W$

$$\sin^2\theta_W^{\text{On-shell}} = 0.2277 \pm 0.0013 \text{ (stat)} \pm 0.0009 \text{ (syst)}$$
$$\Rightarrow M_W^{\text{On-Shell}} = 80.14 \pm 0.08 \text{ GeV}/c^2$$

- NuTeV result deviates from SM prediction by about $+3\sigma$ (PRL 88, 091802, 2002)
- Interpretations of this result implicates lower left-hand coupling (-2.6σ) but good agreement in right-hand coupling with SM
- NuTeV discrepancy has generated a lot of interest in the community
 - Still could be a large statistical fluctuation (5σ has happened before)
 - No single one can explain the discrepancy
- NuTeV working on NLO analysis of $\sin^2\theta_W$
- A Neutrino factory can provide a dramatic improvement in $\sin^2\theta_W$
 - Large neutrino flux (both ν_e and ν_μ)
 - Significant improvement in uncertainties ($\Delta M_W < 25 \text{ MeV}$)
 - Light target detector with p-id would be necessary
 - Theoretical improvement will help further improving the measurement