

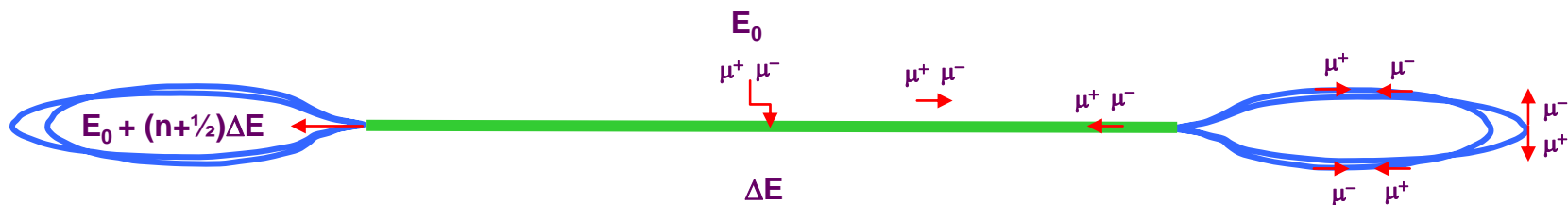
Muon Acceleration in 'Dogbone' RLAs

Alex Bogacz

Overview

- Dogbone configuration
 - orbit separation
 - simultaneous acceleration of both μ^+ μ^- species
- **12.6 GeV Two-step-Dogbone RLA (4.5 pass) – ISS**
- Focusing scheme – Triplet vs FODO lattices
 - multi-pass linac optics
 - phase slippage in the linac
 - 'droplette' Arc lattice
- 15 GeV Dogbone RLA (6.5 pass)

Simultaneous acceleration of both μ^+ μ^- species



orbit separation at linac's end \sim energy difference between consecutive passes ($2\Delta E$)

Baseline Acceleration Scenario (ISS)

- The scheme involves three superconducting linacs (200 MHz, 15 MeV/m):
 - a single pass linear Pre-accelerator
 - followed by a pair of multi-pass 'Dogbone' recirculating linacs (RLAs).
- Acceleration starts after ionization cooling at 273 MeV/c and proceeds to 12.6 GeV/c
- The beam may be injected into FFAG ring(s) for further acceleration

12.6 GeV Two-step-Dogbone RLA

Dogbone I (4.5 pass):

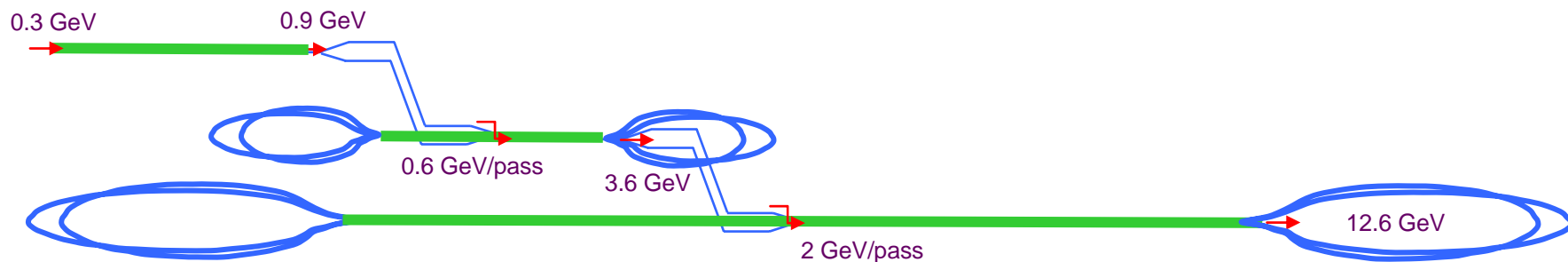
0.9 GeV to 3.6 GeV

$$\frac{E_f}{E_0} = 3.5$$

Dogbone II (4.5 pass):

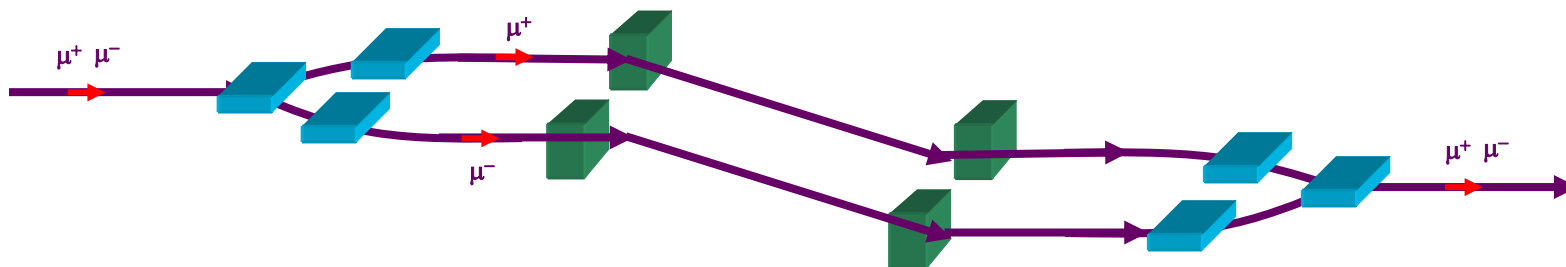
3.6 GeV to 12.6 GeV

$$\frac{E_f}{E_0} = 4$$

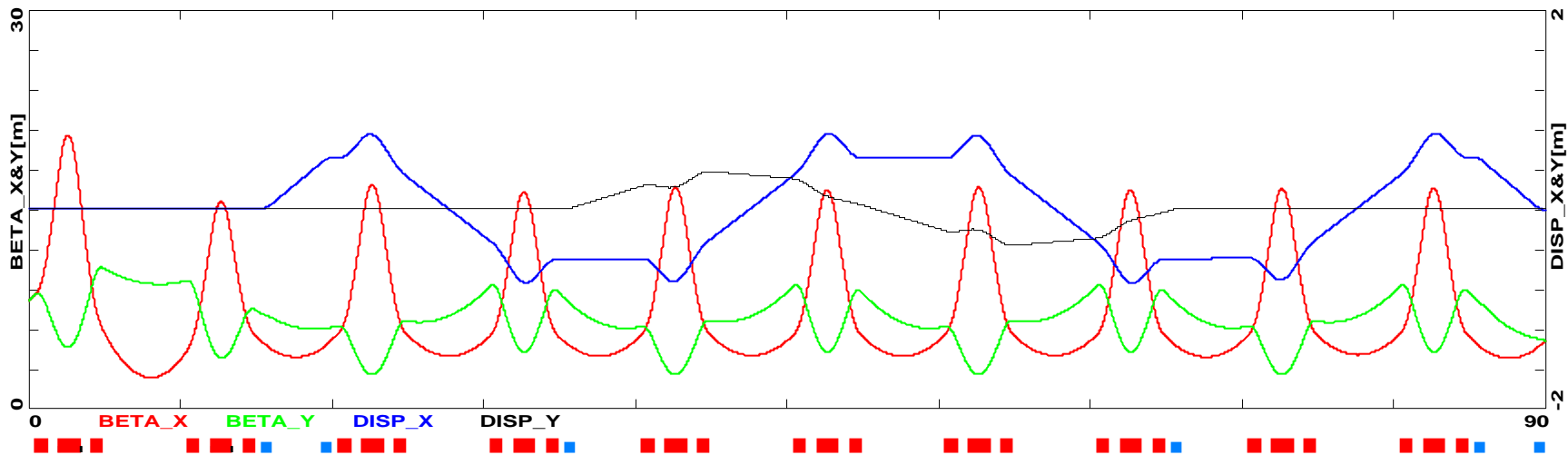


For compactness all these components (Pre-accelerator, Dogbone I and Dogbone II) are stacked up vertically; μ^\pm beam transfer between the accelerator components is facilitated by the vertical double chicane

Injection double-chicane



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Phase slippage in the linac

- The energy range for each 'Dogbone' RLA was chosen to give similar ratios of top-to-injection energies:
 - 3.5 for Dogbone I
 - 4.0 for Dogbone II
- The phase slippage of a muon injected with the initial energy E_0 and accelerated by ΔE in a linac of length, L , was calculated, where uniformly spaced RF cavities were phased for a speed-of-light particle
- The injection energies for both RLAs were chosen, so that a tolerable level of the RF phases slippage along a given length linac can be maintained.

Phase slippage in the linac

- A simple calculation of the phase slippage of a semi-relativistic muon accelerated in a linac, where uniformly spaced RF cavities are phased for the speed-of-light particle was carried out using the following cavity-to-cavity iterative algorithm for phase-energy vector

$$\begin{pmatrix} \phi_{k,i+1} \\ E_{k,i+1} \end{pmatrix} := \begin{bmatrix} \phi_{k,i} + \frac{h}{\lambda} \cdot 360 \left[\frac{1}{2} \cdot \left(\frac{m_{\mu}}{E_{k,i}} \right)^2 \right] \\ E_{k,i} + h \cdot \Delta E_k \end{bmatrix}$$

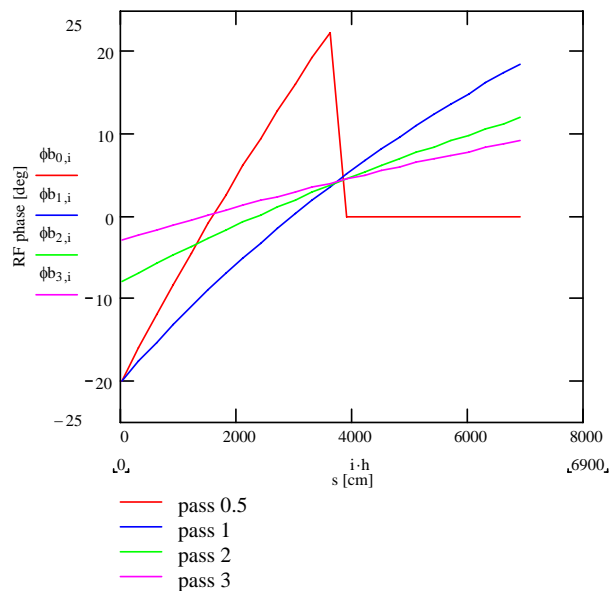
where

$$h := \frac{L_{\text{linac}}}{N_{\text{cav}}} \quad \lambda := \frac{c}{f_0} \quad k := 0..4 \quad i := 0..1N_{\text{cav}} - 1$$

Phase slippage in the linacs

Dogbone I

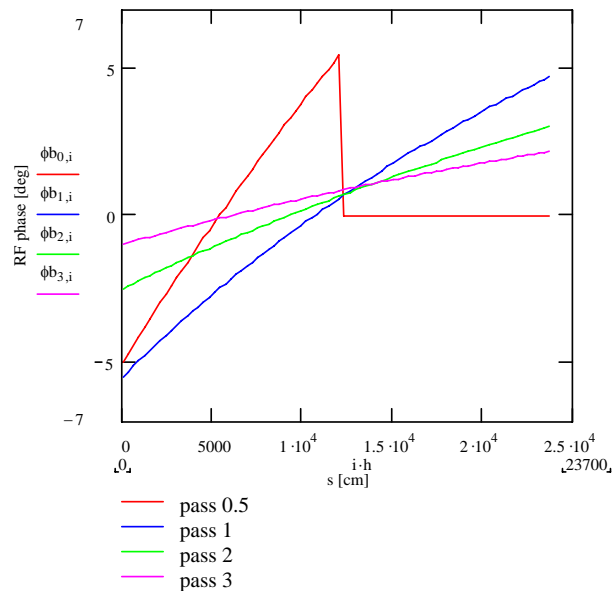
$$E_0 = 0.9\text{GeV}, \Delta E = 0.6\text{GeV}, L = 70\text{m}$$



$$\begin{aligned} \phi_{0,0} &:= -20 & \phi_{1,0} &:= -20 \\ \phi_{2,0} &:= -8 & \phi_{3,0} &:= -3 \end{aligned}$$

Dogbone II

$$E_0 = 3.6\text{GeV}, \Delta E = 2\text{GeV}, L = 240\text{m}$$



$$\begin{aligned} \phi_{0,0} &:= -5 & \phi_{1,0} &:= -5.5 \\ \phi_{2,0} &:= -2.5 & \phi_{3,0} &:= -1 \end{aligned}$$

RF phase slippage along the multi-pass linacs; initial 'gang phases' for each pass were chosen for the optimum longitudinal bunch compression in each linac-Arc segment

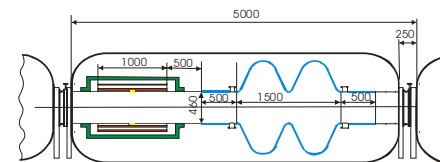
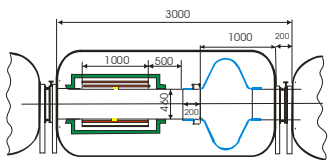
Longitudinal compression in the RLA

- Longitudinal bunch compression is required in the course of acceleration.
- To accomplish that, the beam is accelerated off-crest with non zero $M_{56} \sim 6$ m (momentum compaction) in the 'droplette' Arcs.
- This induces synchrotron motion, which suppresses the longitudinal emittance growth related to non-linearity of accelerating voltage.
- Without synchrotron motion the minimum beam energy spread would be determined by non-linearity of RF voltage across the bunch length, e.g. it would be equal to $1 - \cos\phi \approx 9\%$ for bunch length $\phi = 30$ deg.
- The synchrotron motion within the bunch averages the total energy gain of tail's particle to the energy gain of particles in the core.

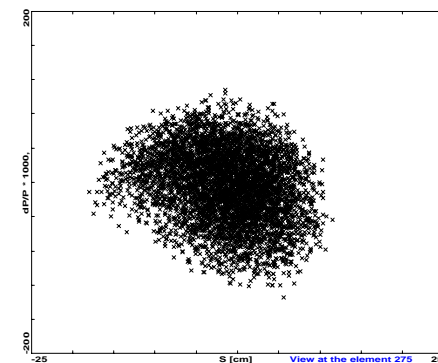
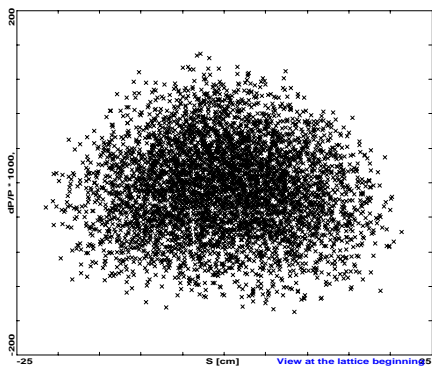
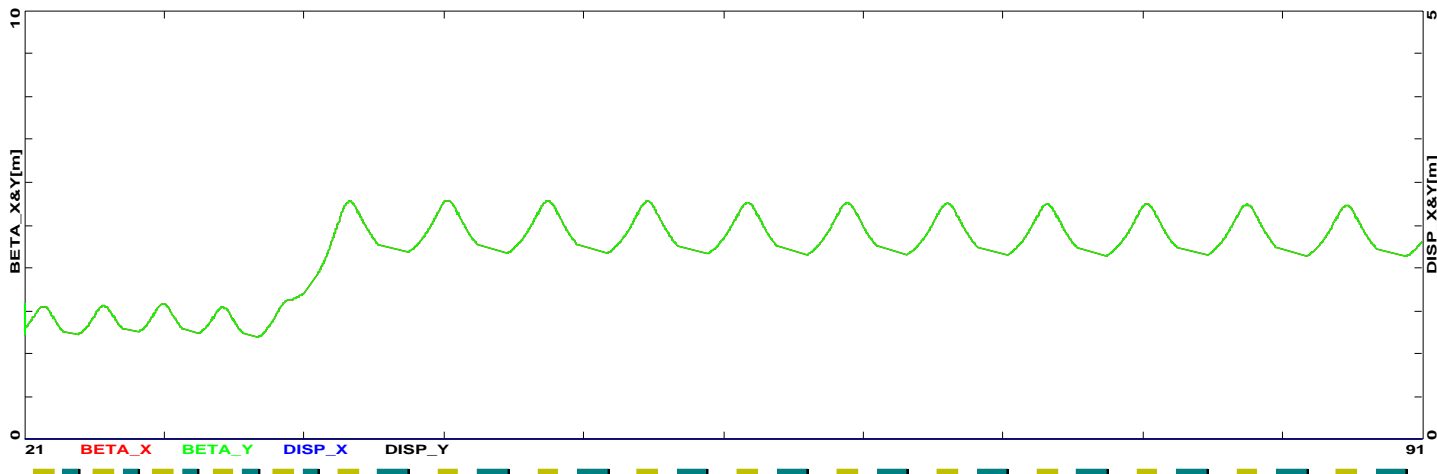
Initial beam emittance/acceptance after cooling at 273 MeV/c

Normalized Emittances		ϵ_{rms}	$A = (2.5)^2 \epsilon$
transverse emittance: ϵ_x/ϵ_y	mm·rad	4.8	30
longitudinal emittance: ϵ_l ($\epsilon_l = \sigma_{\Delta p} \sigma_z / m_\mu c$)	mm	27	150
momentum spread: $\sigma_{\Delta p/p}$		0.07	± 0.17
bunch length: σ_z	mm	176	± 442

Linear Pre-accelerator – Longitudinal dynamics

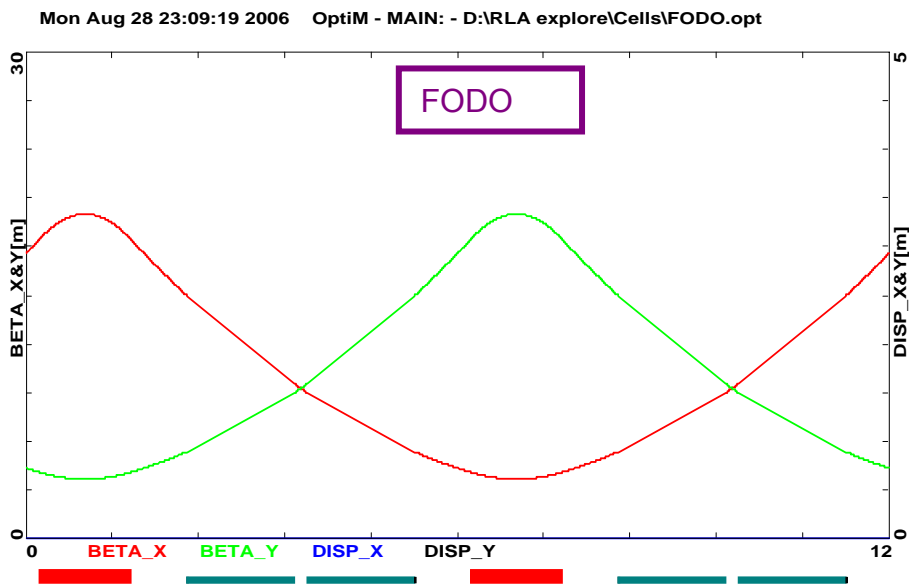


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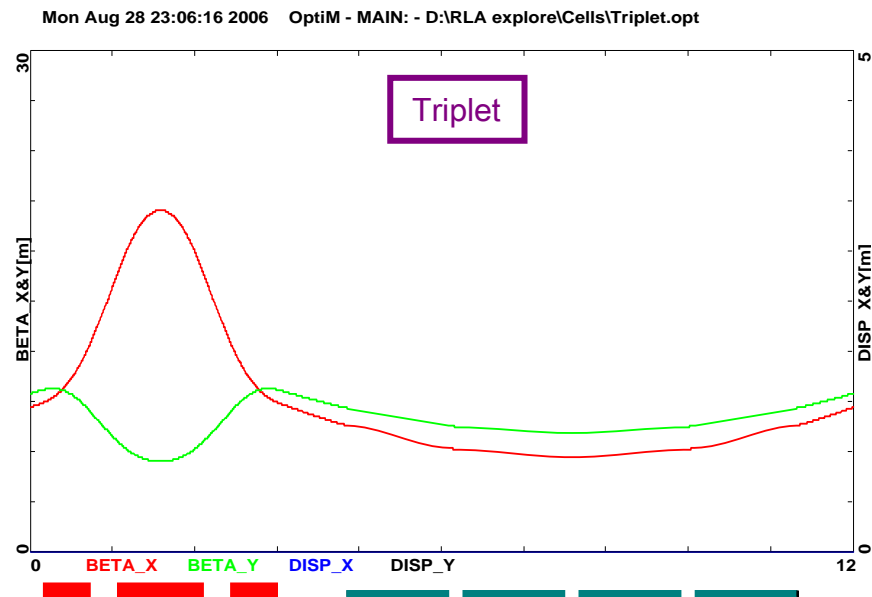
FODO vs Triplet focusing structure

- The same length
- The same phase advance per cell ($\Delta\phi_x = 90^\circ = \Delta\phi_y$)



Advantages:

- much weaker quads (~3 times)
- shorter quads (total)
- easier chromaticity correction



Advantages:

- longer straight sections
- smaller vertical beta-function
- uniform variation of betas and dispersion

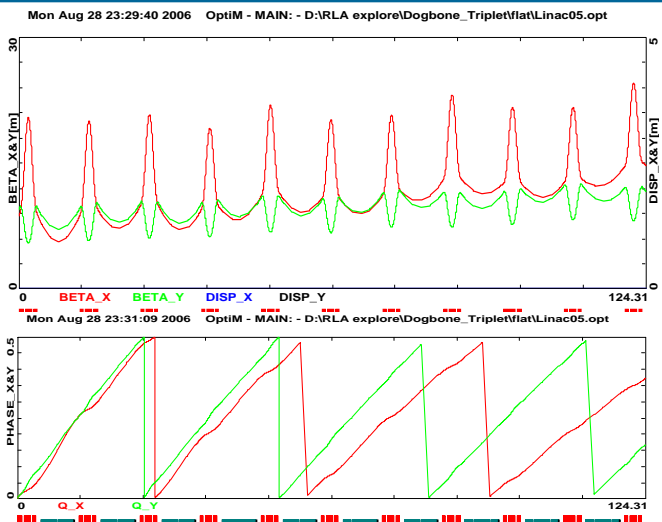
Multi-pass Linac Optics

- The focusing profile along the linac (quadrupole gradients) need to be set so that one can transport multiple pass beams within a vast energy range (provide adequate transverse focusing for given aperture) .
- The beam is traversing the linac in both directions – one chooses a ‘flat focusing profile’ (Bob Palmer) for the entire linac: e.g. the quads in all cells are set to the same gradient, corresponding to 90 deg. phase advance per cell determined for the lowest energy (injection) – no quad scaling with energy
- The requirement of simultaneous acceleration of both μ^\pm species imposes mirror symmetry of the ‘droplette’ Arcs optics (the two species move in the opposite directions through the Arcs). This in turn puts a constraint on the exit/entrance Twiss functions for the two consecutive linac passes:

$$\beta_n^{\text{out}} = \beta_{n+1}^{\text{in}} \text{ and } \alpha_n^{\text{out}} = -\alpha_{n+1}^{\text{in}}$$

where $n = 0, 1, 2..$ is the pass index

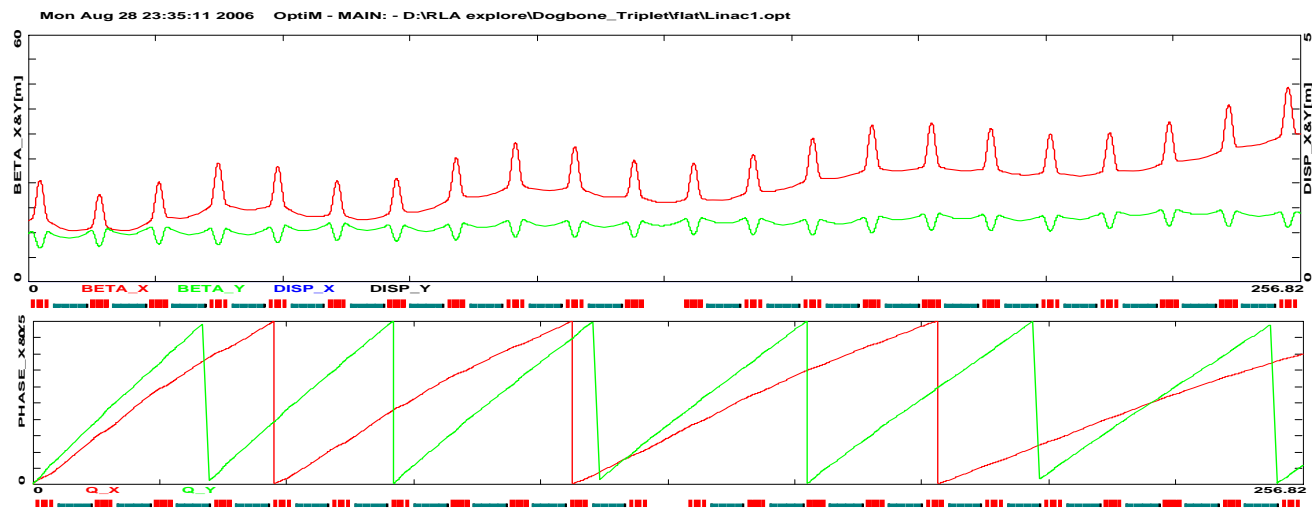
mirror symmetry cond. ($\beta_{out}^n = \beta_{in}^{n+1}$, and $\alpha_{out}^n = -\alpha_{in}^{n+1}$, n - pass index)



'half pass' , 2-3 GeV



Initial phase adv/cell 90 deg – fixed quad gradients for the entire linac (no scaling with energy)



1-pass, 3-5 GeV

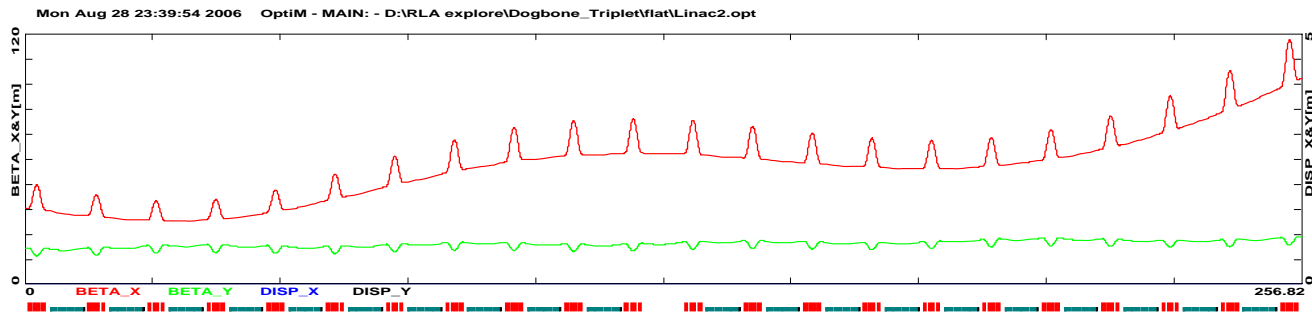


phase adv. drops much faster in the horizontal plane

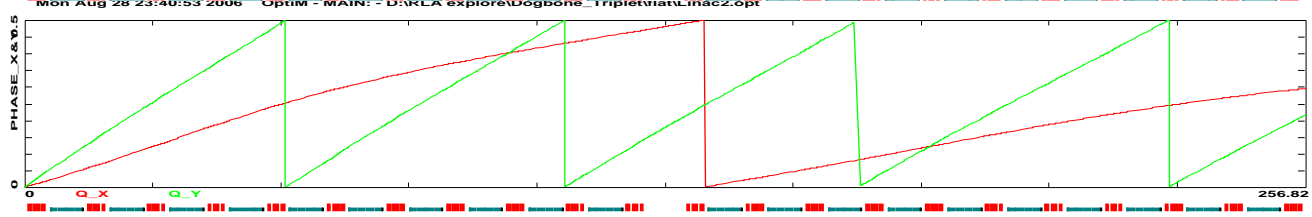
Triplet - 'flat focusing' linac profile



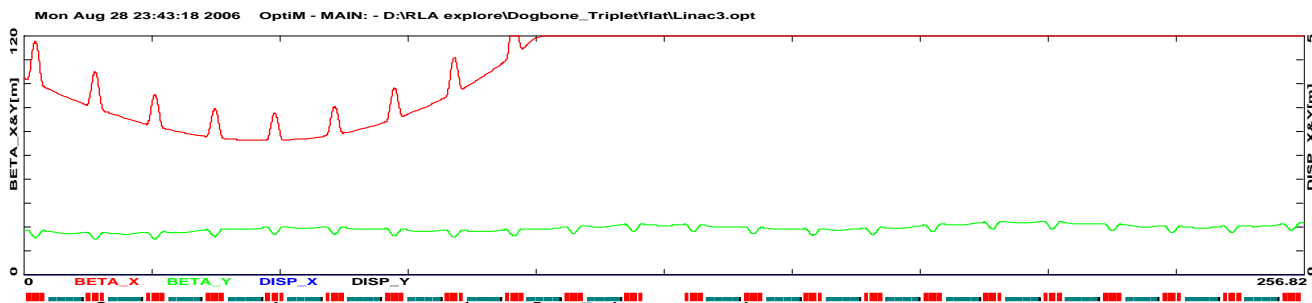
mirror symmetry cond. ($\beta_{out}^n = \beta_{in}^{n+1}$, and $\alpha_{out}^n = -\alpha_{in}^{n+1}$, n - pass index)



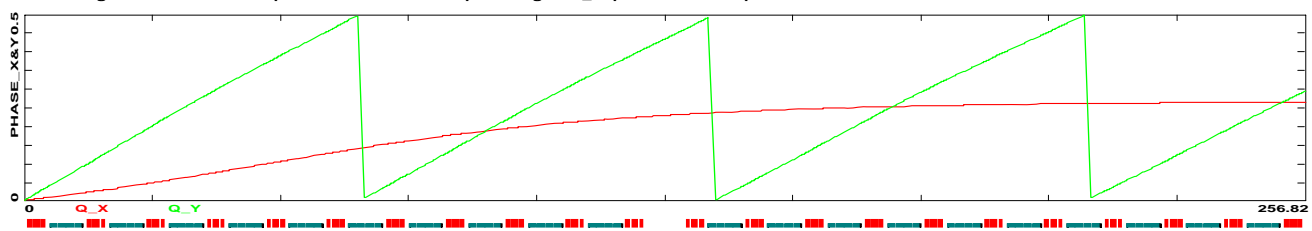
2-pass, 5-7 GeV



phase adv. drops much faster in the horizontal plane



3-pass, 7-9 GeV



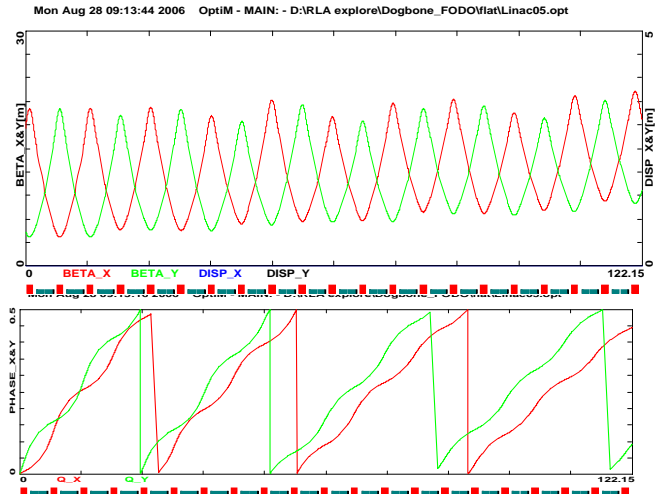
no phase adv. in the horizontal plane



FODO - 'flat focusing' linac profile



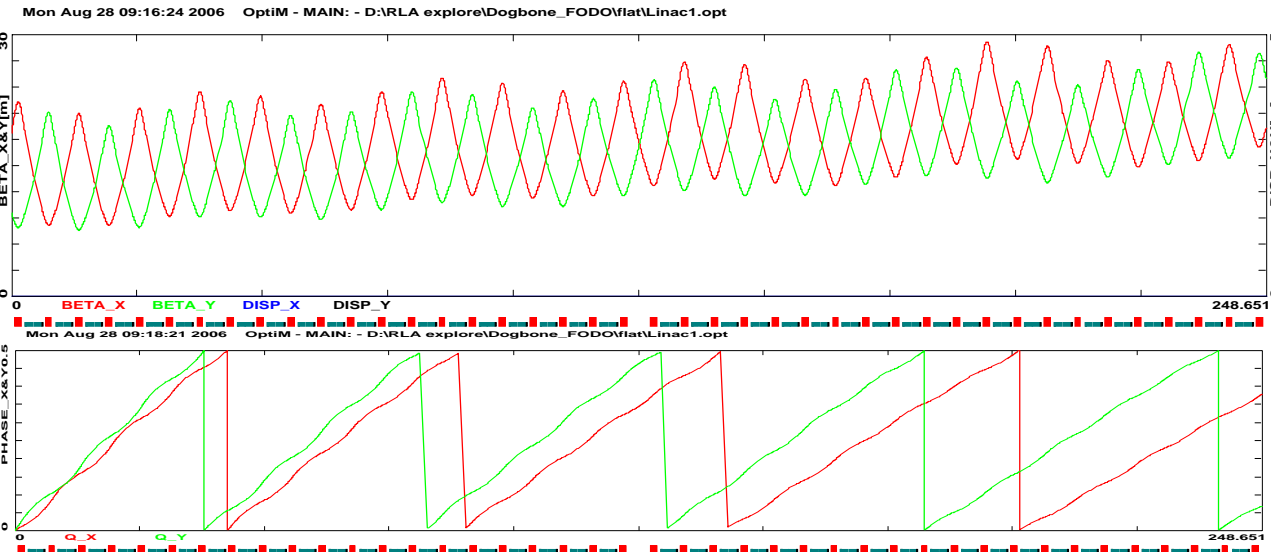
mirror symmetry cond. ($\beta_{out}^n = \beta_{in}^{n+1}$, and $\alpha_{out}^n = -\alpha_{in}^{n+1}$, n - pass index)



'half pass', 2-3 GeV



initial phase adv/cell 90 deg – fixed gradient in all cells
(no scaling with energy)



1-pass, 3-5 GeV



phase adv. diminish uniformly in both planes

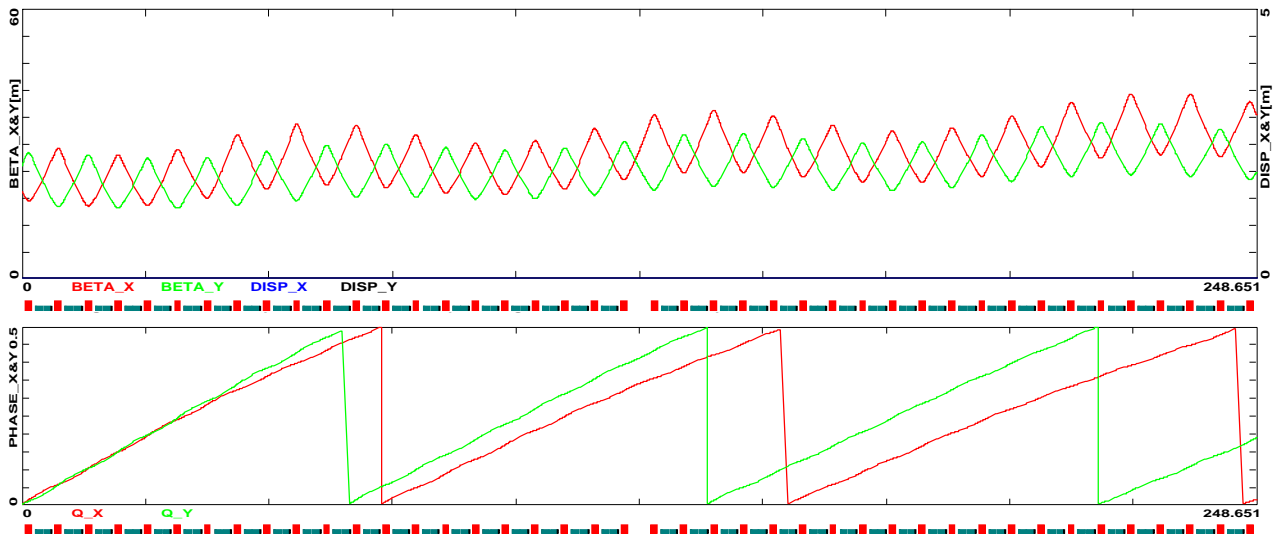


FODO - 'flat focusing' linac profile



mirror symmetry cond. ($\beta_{out}^n = \beta_{in}^{n+1}$, and $\alpha_{out}^n = -\alpha_{in}^{n+1}$, n - pass index)

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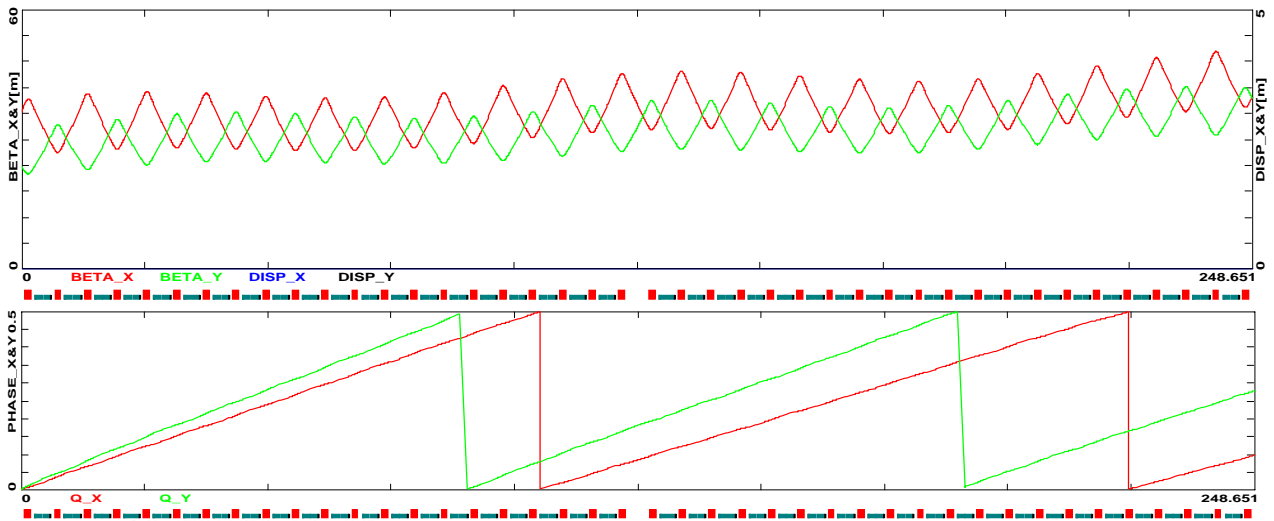


2-pass, 5-7 GeV



phase adv. diminish uniformly in both planes

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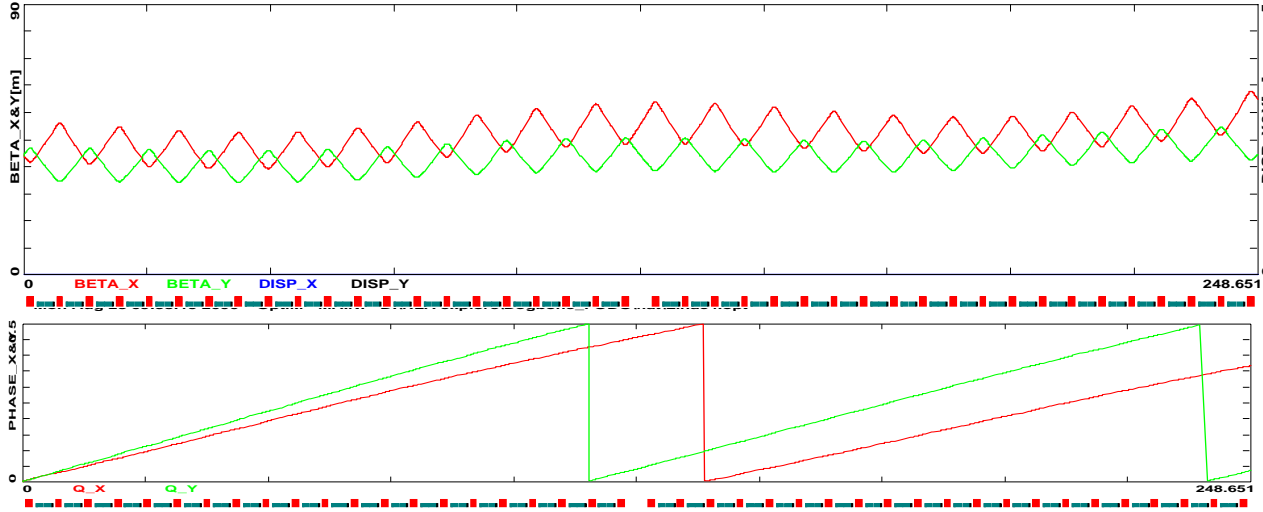
3-pass, 7-9 GeV



FODO - 'flat focusing' linac profile

mirror symmetry cond. ($\beta_{out}^n = \beta_{in}^{n+1}$, and $\alpha_{out}^n = -\alpha_{in}^{n+1}$, n - pass index)

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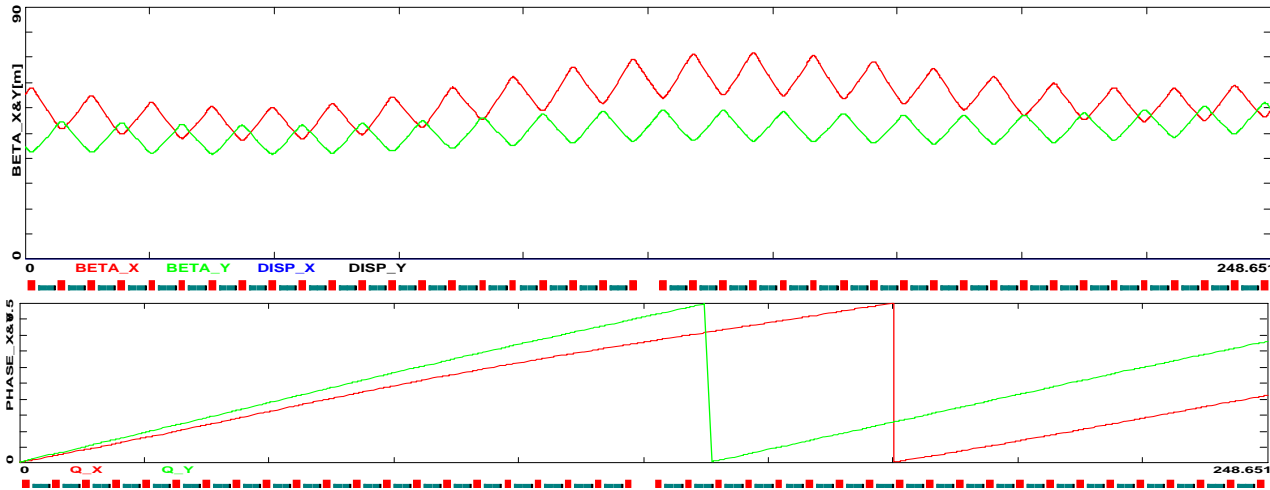


4-pass, 9-11 GeV



phase adv. diminish uniformly in both planes

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5-pass, 11-13 GeV

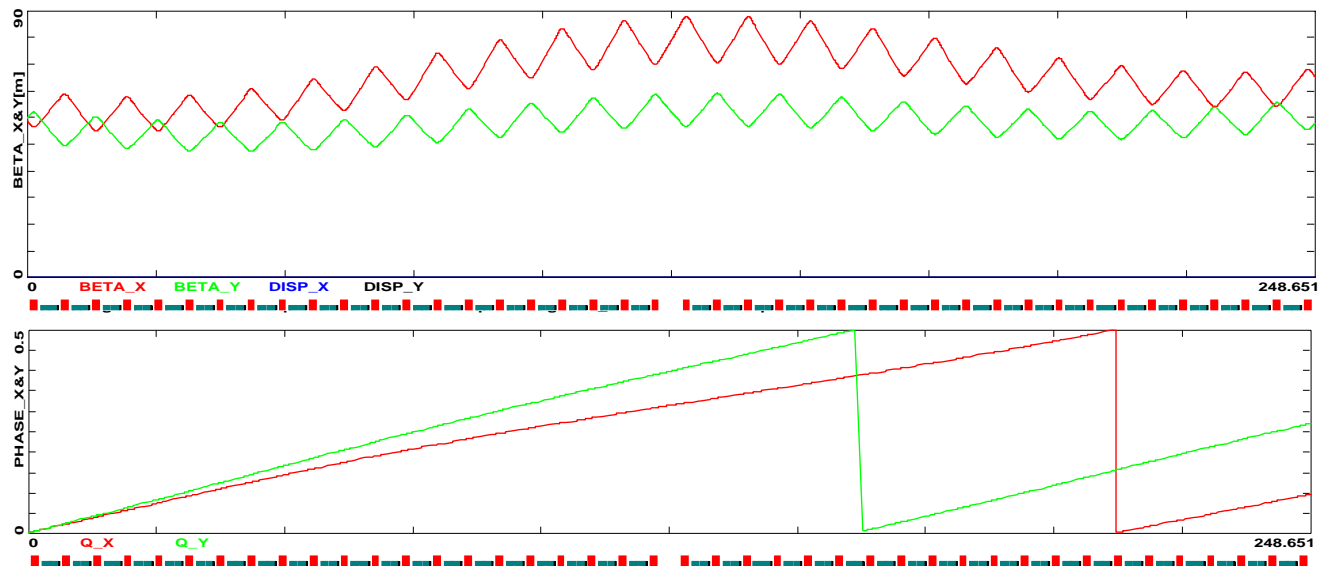


'flat focusing' linac profile



mirror symmetry cond. ($\beta_{out}^n = \beta_{in}^{n+1}$, and $\alpha_{out}^n = -\alpha_{in}^{n+1}$, n - pass index)

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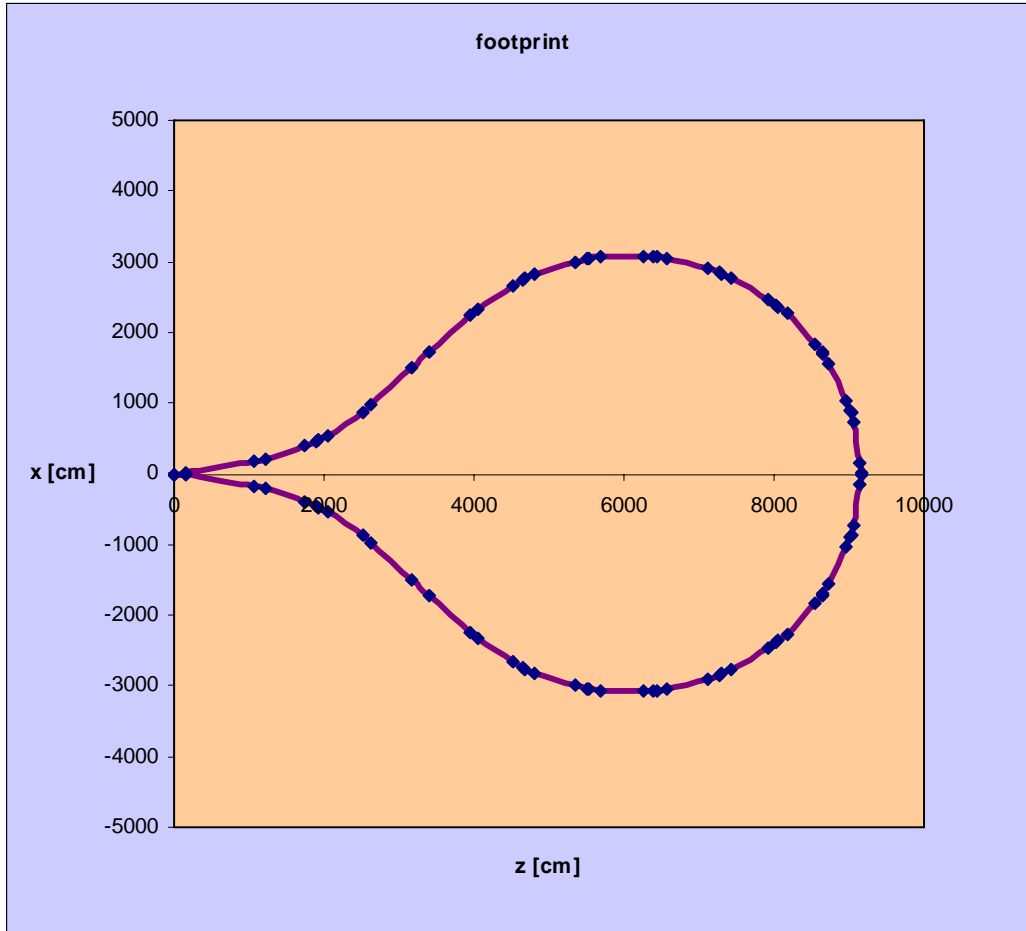
6-pass, 13-15 GeV



phase adv. diminish uniformly in both planes



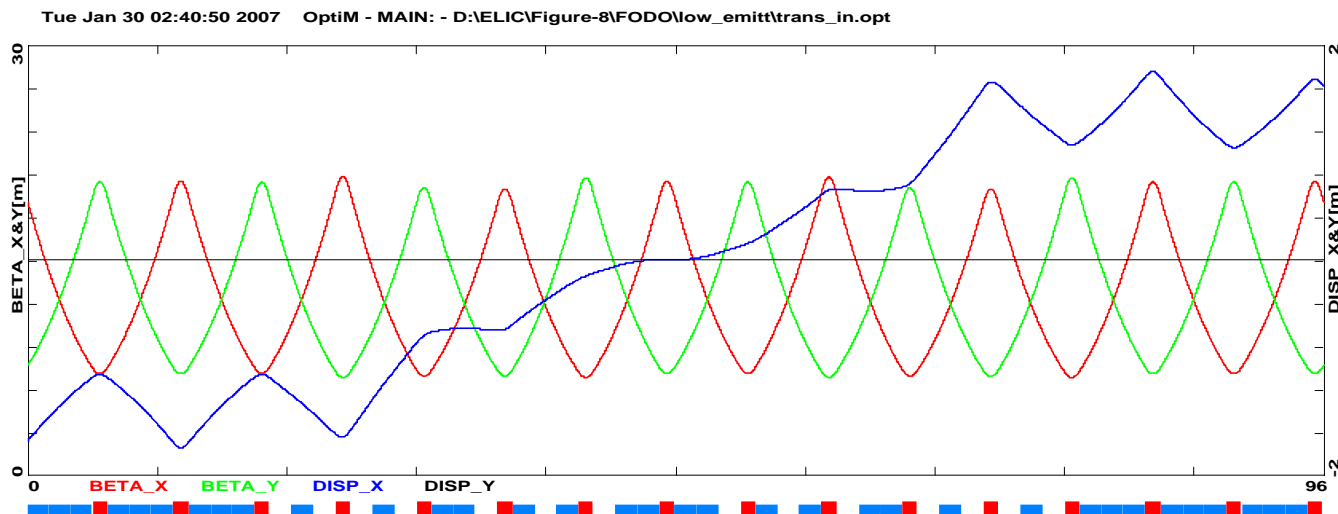
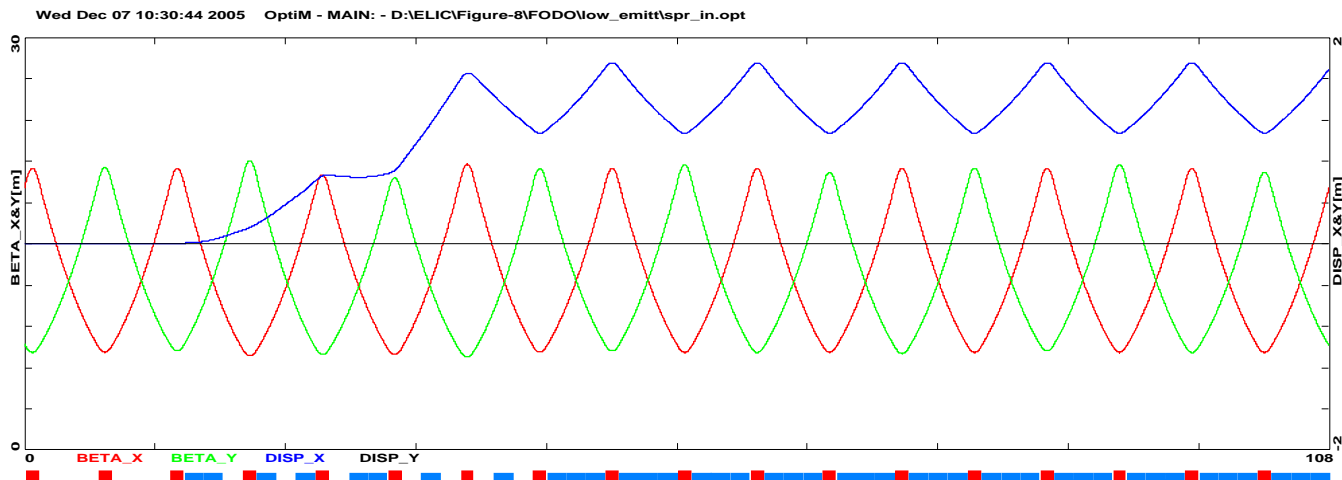
'Droplette' Arc – Layout



Arc dipoles

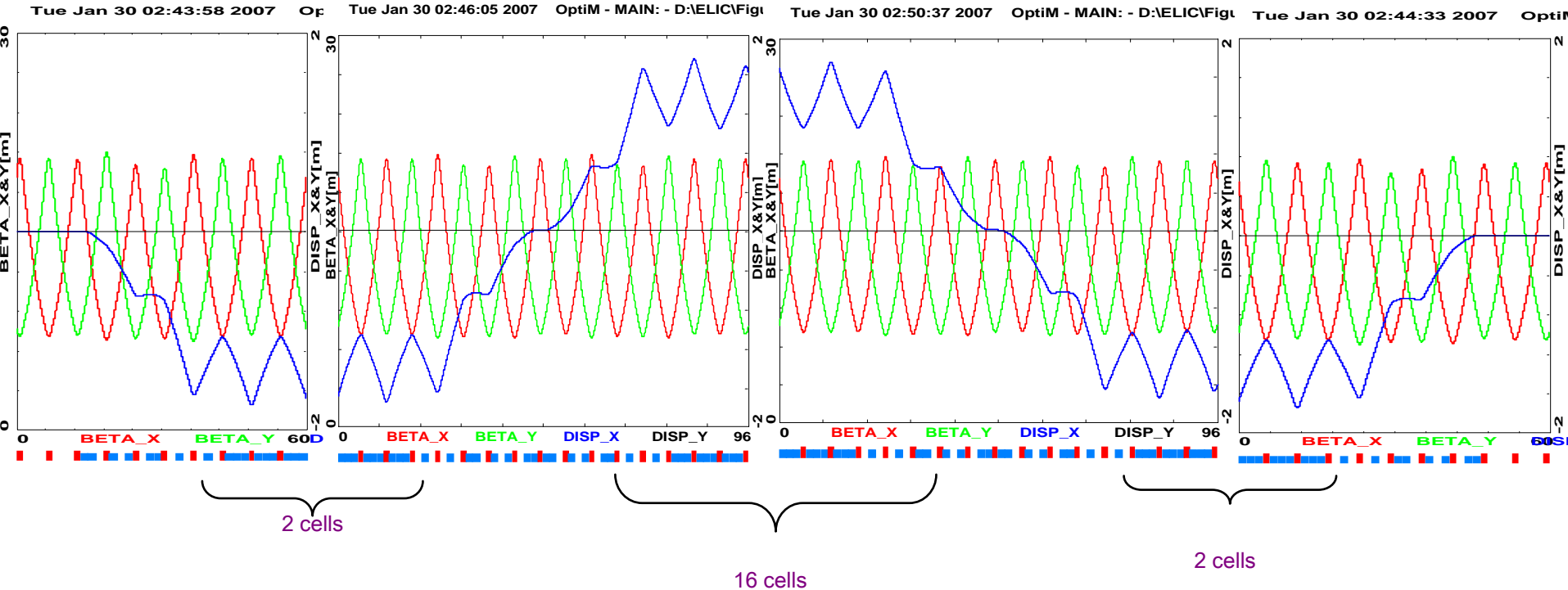
$L_b=150$; \Rightarrow 150 cm
 $\theta_0=10.3283$; \Rightarrow 10.328 deg
 $N_{in}=16$; \Rightarrow 16
 $N_{out}=2$; \Rightarrow 2
 $\theta = (90 + \theta_0) / (N_{in} - 2 * N_{out})$; \Rightarrow 8.36 deg.
 #
 $\theta_{out} = \theta_0 + 2 * N_{out} * \theta$; \Rightarrow 43.77 deg.
 $\theta_{in} = 2 * N_{in} * \theta$; \Rightarrow 267.54 deg.
 $B_p = \pi * H_r * \theta / (180 * L_b)$; \Rightarrow 6.537 kGauss
 $L_{ring} = 227.3$ m

Spreader and 'Dispersion Flip' Lattices



'Droplette' Arc – Mirror-symmetric Optics

$(\beta_{out} = \beta_{in}, \text{ and } \alpha_{out} = -\alpha_{in}, \text{ matched to the linacs})$



$$M_{56} = - \int \frac{D_x}{\rho} ds = -D_x^{dip} \int d\left(\frac{s}{\rho}\right) = -D_x^{dip} \int d\theta_{rad} = -D_x^{dip} \times \theta_{rad}^{tot}$$

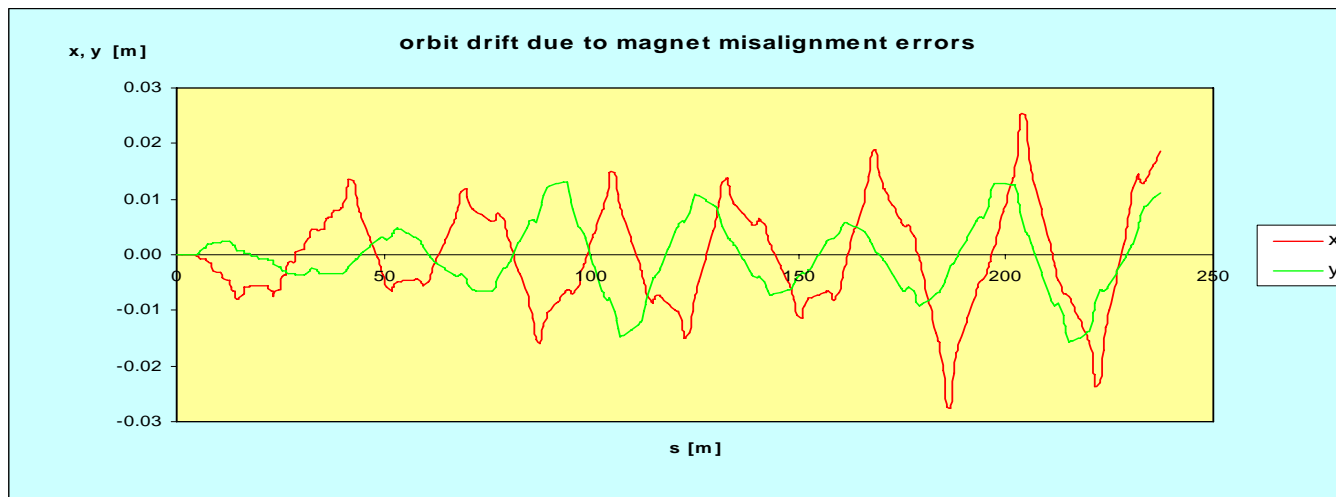
$$M_{56} = -\frac{5}{3} \pi \times 1.2 \text{ m} = -6.3 \text{ m}$$

Magnet Misalignment Errors

- Lattice sensitivity to random misalignment errors was studied via DIMAD Monte-Carlo assuming:
quadrupole misalignment errors:

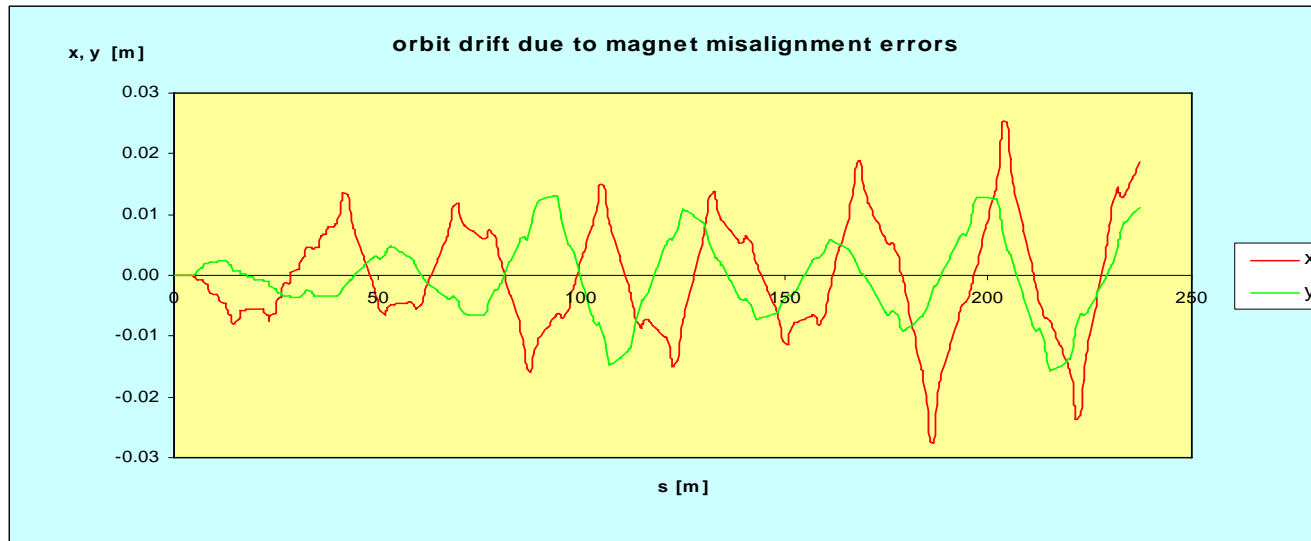
F:	$\sigma_x = \sigma_y = 1 \text{ mm}$	$(\sigma_{x,y'} = \sigma_{x,y}/L)$	$\left\{ \begin{array}{l} \sigma_{x'} = \sigma_{y'} = 0.8 \times 10^{-3} \\ \sigma_{x''} = \sigma_{y''} = 1.47 \times 10^{-3} \end{array} \right.$
D:	$\sigma_x = \sigma_y = 1 \text{ mm}$		

- Gaussian distribution was chosen for individual quad misalignments
- Resulting reference orbit distortion (uncorrected) for Arc 2 is illustrated below



- Similar level of dipole misalignment errors had virtually no effect on random steering

Arc 2 – Magnet Misalignment Errors



RMS Orbit Displacement [m]:	
X:	0.9486e-02
y:	0.7003e-02

Extr. Orbit Displacement [m]:	
X_{max} :	0.2538E-01
X_{min} :	-0.2782E-01
y_{max} :	0.1434E-01
y_{min} :	-0.1697E-01

- Same level of orbit drifts due to quad misalignments for other ‘Dogbone’ segments (Arc 1, 3 and 4 and linacs)
- Orbit drifts at the level of ~3 cm can easily be corrected by pairs of hor/vert correctors (2000 Gauss cm each) placed at every triplet girder

Focusing Error Tolerances – Quadrupole Field Spec

- Cumulative Arc-to-Arc Optics mismatch as measured by Courant-Snyder invariant change:

$$\begin{aligned}\varepsilon' &= \beta(\theta + \delta\theta)^2 + 2\alpha(\theta + \delta\theta)x + \gamma x^2 \\ &= \varepsilon \left(1 + \beta\Delta\Phi \sin(2\mu) + (\beta\Delta\Phi \cos \mu)^2 \right),\end{aligned}$$

$$\begin{aligned}\varepsilon_N &= \varepsilon_0 \prod_{n=1}^N \left(1 + \beta_n \Delta\Phi_n \sin(2\mu_n) + (\beta_n \Delta\Phi_n \cos \mu_n)^2 \right) \\ &= \varepsilon_0 \prod_{n=1}^N \left(1 + \frac{1}{2} (\beta_n \Delta\Phi_n)^2 + \sqrt{(\beta_n \Delta\Phi_n)^2 + \left(\frac{\beta_n \Delta\Phi_n}{2} \right)^4} \sin(2\mu_n + \psi_n) \right),\end{aligned}$$

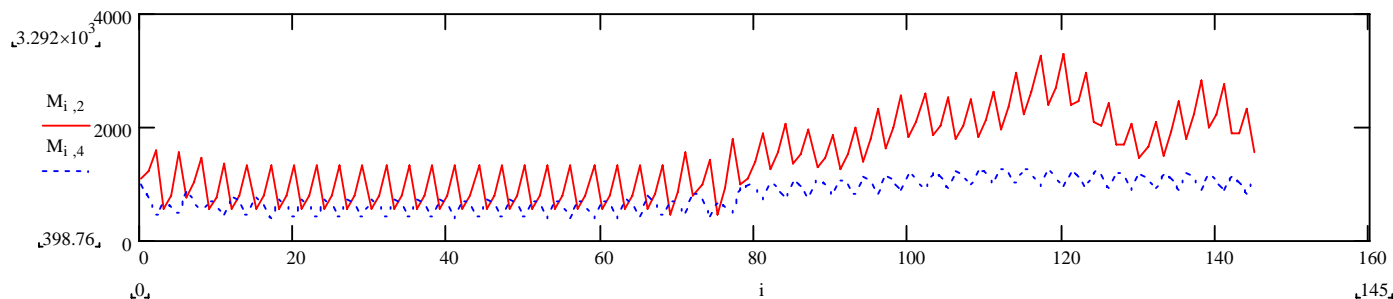
- Standard deviation of Courant-Snyder invariant:

$$\frac{\sigma_\varepsilon}{\varepsilon} = \frac{\sqrt{\Delta\varepsilon^2 - \overline{\Delta\varepsilon^2}}}{\varepsilon} \approx \sqrt{\frac{1}{2} \sum_{n=1}^N (\beta_n \Delta\Phi_n)^2} = \frac{\sqrt{\Delta\Phi^2}}{\Phi_{\max}} \sqrt{\frac{1}{2F_{\min}^2} \sum_{n=1}^N \beta_n^2}$$

Focusing Error Tolerances – Quadrupole Field Spec

- By design, one can tolerate Arc-to-Arc mismatch at the level of **10%** (to be compensated by the dedicated matching quads).

$$\left(\frac{\sigma_\varepsilon}{\varepsilon}\right)_{mis} = \sqrt{\frac{1}{2} \sum_{n=1}^N (\beta_n \delta\phi_1)^2} = \sqrt{\frac{1}{2} \Delta\phi_1^2 \sum_{n=1}^N (\beta_n)_{quad}^2 + \frac{1}{2} \delta\phi_1^2 \sum_{n=1}^N (\beta_n)_{dipole}^2}$$



(GdL=65 kG)

$F_{min} = 1.5 \text{ m}$

$$\sqrt{\frac{1}{2F_{min}^2} \sum_{n=1}^N \beta_n^2} \approx 190$$

$$\frac{\Phi}{\phi_{max}} = 0.001$$

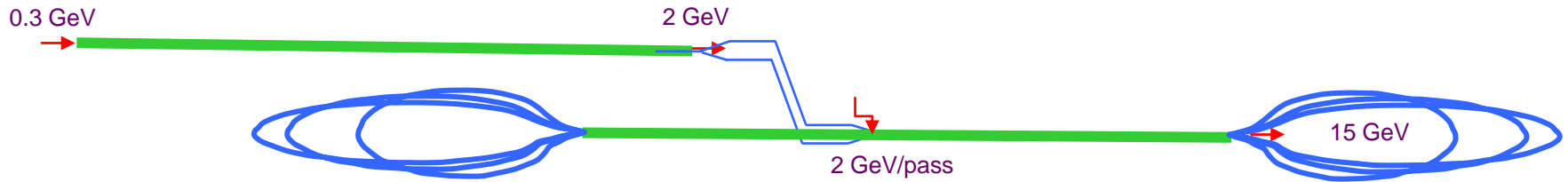
$$\Phi = \sqrt{\delta\phi_1^2 + \frac{3}{2} a^2 (\delta\phi_2^2 + 2\delta\phi_1\delta\phi_3) + \frac{5}{2} a^4 (\delta\phi_3^2 + 2\delta\phi_1\delta\phi_5 + 2\delta\phi_2\delta\phi_4) + \dots}$$



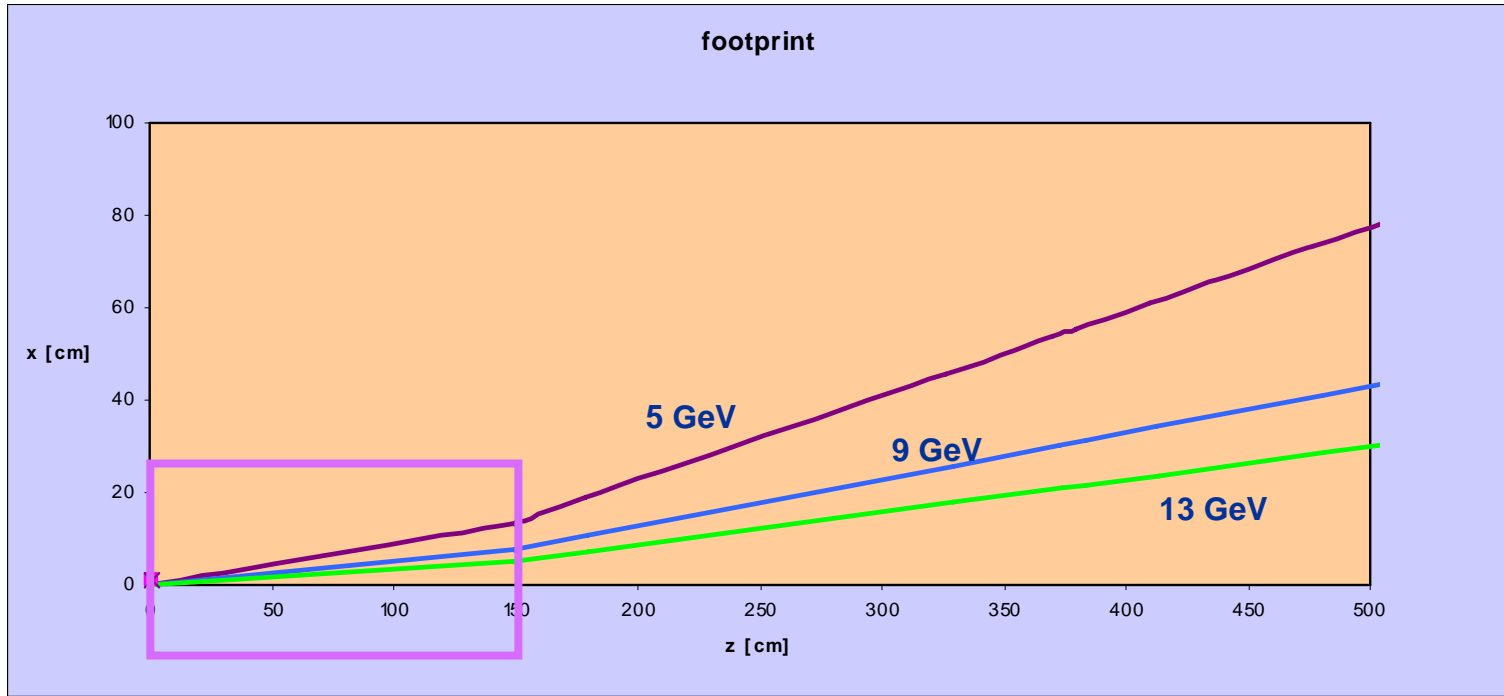
15 GeV Dogbone RLA (6.5 pass)

energy ratio:

$$\frac{E_f}{E_0} = 7.5$$



Beam Separation – Switchyard



Summary

- **'Dogbone' RLA** – preferred configuration
 - better orbit separation for higher passes
 - offers symmetric solution for simultaneous acceleration of μ^+ and μ^-
- **FODO lattice** more favorable (compared to the triplet) to accommodate large number of passes
 - uniform phase advance decrease in both planes
 - smaller variation of Twiss function – easier match to the Arcs
- **Proposed 12.6 GeV two-step-dogbone RLA** (30 mm rad acceptance)
- **3.6 GeV Dogbone I** (4.5 pass) – error sensitivity studies
 - Magnet misalignment error analysis (DIMAD Monte Carlo on the above lattice) shows quite manageable level of orbit distortion for ~ 1 mm level of magnet misalignment error).
 - Great focusing errors tolerance for the presented lattice – 10% of Arc-to-Arc betatron mismatch limit sets the quadrupole field spec at 0.1%
- **Aggressive 15 GeV dogbone RLA**