# An Induction Kicker for Muon Cooling Rings

and other Applications Needing Very Large Apertures

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Study 2 cooling is 100 m If Ring is to be cheaper: Circ ≪ 100 m

say 30 m

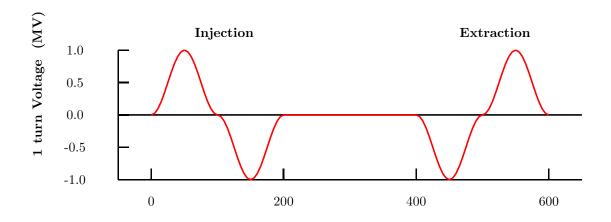
e.g. Balbakov 4 sided ring, Palmer 11 cell ring

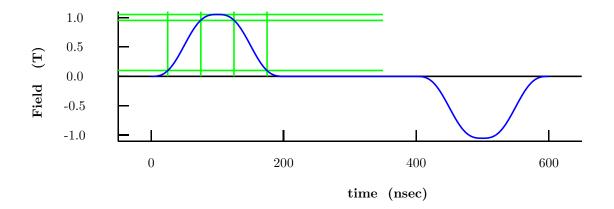
Balbakov 8 cell (60m) more expensive Palmer 5 cell (16 m circ) req faster kicker 30m is Compromise

Time = 100 nsec

50 nsec for train
50 nsec for kicker rise

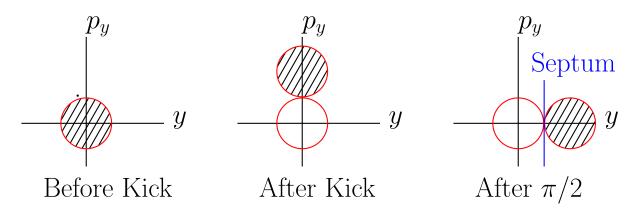
# Required Volts and Current





- If req.  $\delta B \leq +/-5 \%$
- circ/2 = rise 10% to 90%=  $50 \operatorname{nsec}$

# Minimum Required kick



# **Defining**

For Aperture:

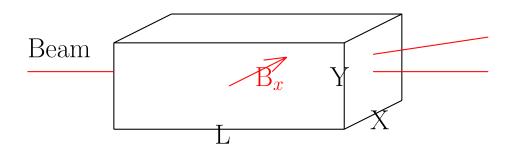
$$f_{\sigma} = \frac{\mathrm{Ap}}{\sigma}$$

For leakage flux:

$$\Phi = f_{\Phi} B_x L Y$$

For finite  $\mu$ 's:

$$\int Bdl/\mu = f_{\mu} B X$$



#### **Formulae**

$$\Delta p_y = B_x L c = m_\mu 2 f_\sigma \sqrt{\frac{\epsilon_n \beta_\gamma}{\beta_y}}$$

$$X = 2 f_\sigma \sqrt{\frac{\epsilon_n \beta_x}{\beta_\gamma}}$$

$$R = \frac{Y}{X} \approx \left(1 + \frac{L}{2 \beta_y}\right)$$

$$I = \frac{f_{\mu} B X}{\mu_o} = \frac{f_{\mu} 4 f_{\sigma}^2 m_{\mu}}{\mu_o c} \frac{\epsilon_n}{L}$$

$$V = \frac{f_{\Phi} B Y L}{t_{\text{rise}}} = \frac{f_{\Phi} 4 f_{\sigma}^2 m_{\mu} R}{c} \frac{\epsilon_n}{\tau}$$

$$J = f_{\mu} f_{\Phi} \frac{B^2 L X Y}{2 \mu_o} = f_{\mu} f_{\Phi} \frac{m_{\mu}^2 8 f_{\sigma}^4 R}{\mu_o c^2} \frac{\epsilon_n^2}{L}$$

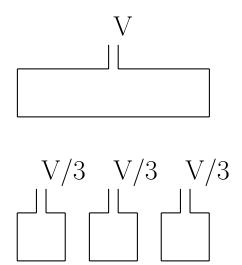
#### Note

- V independent of E,  $\beta_{\perp}$
- V & I  $\propto \epsilon_n$
- J  $\propto \epsilon_n^2$
- muon  $\epsilon_n \gg$  other  $\epsilon_n$ 's
- So muon kicker Joules ≫ other kickers
- Nearest are  $\bar{p}$  kickers

		$\mu$ Cooling	CERN $\bar{p}$	Ind Linac
$egin{array}{c} f_{\Phi} \ f_{\mu} \ . \ f_{\sigma} \end{array}$		1.05		
$ f_{\mu} $ .		1.05		
$ f_{\sigma} $		3		
$\mid m_{\mu} \mid$	V	$1.05 \ 10^8$		
c	m/s	$3 \ 10^8$		
$ \epsilon_n $	$\pi \text{ mm}$	10		
$\beta_x$	m	1.0		
s $Bd\ell$	Tm	.30	.088	
L	m	1.0	$\approx 5$	5.0
$t_{ m rise}$	ns	50	90	40
$t_{ m pulse}$	ns	100	500	100
$ar{eta\gamma}$		2		
В	Τ	.30	$pprox\!0.018$	0.6
X	m	.42	.08	
Y	m	.63	.25	
I	kA	105		
$V_{1turn}$	kV	$3,\!970$	800	$5,\!000$
$U_{\rm magnetic}$	J	$10,\!450$	$pprox\!13$	8000
$n_{\mathrm{parallel}}$		20	10	32
$n_{units}$		12	10	10
$I_{\mathrm{p.s.}}$	kA	192		73
$V_{p.s.}$	kV	182	80	190
$U_{\mathrm{p.s.}}$	kJ	870		800

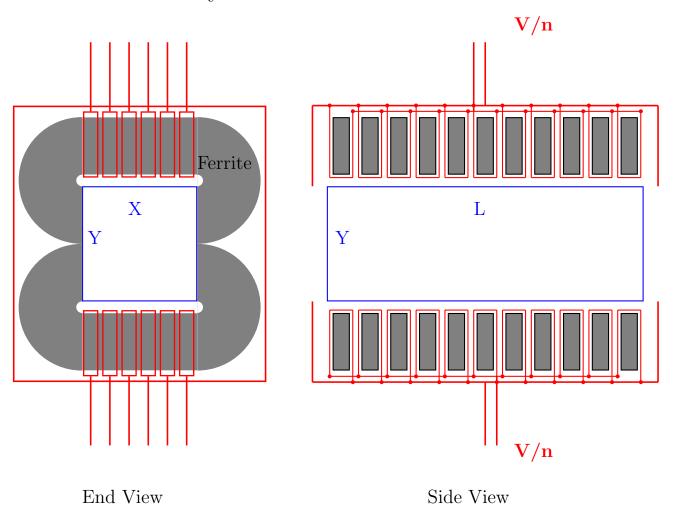
#### Note

- J is 3 orders above  $\bar{p}$
- But same order as 5 m Induction
- And t same order as Induction
- But V is too High
- Because L/X only  $\approx 3$ Conventional subdivision only /3



#### New Idea

- Drive Flux Return
- Subdivide Flux Return
- Conducting Box Removes Stray Field Return

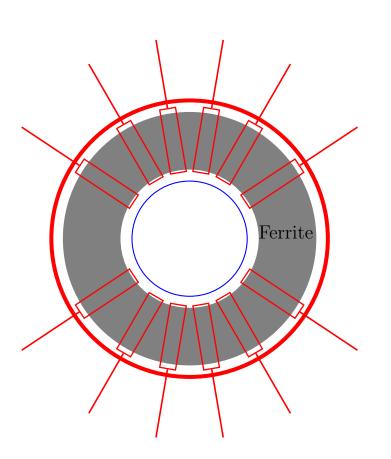


- Divide yoke by  $(2 \times 11)$
- connect 11 loops in parallel
- V / loop = 3970/22 = 182 kV
- I / side =  $105 \times 11 = 1155 \text{ kA}$
- 6 Drivers per yoke, then:
- number of drivers =  $2 \times 6 = 12$
- I / driver = 1155/6 = 192 kA
- U / driver = 1045/12 = 870 J
- U / driver  $\approx$  U / Induction driver
- $\bullet$  V / driver  $\approx$  V / Induction driver
- I is higher from inductive load

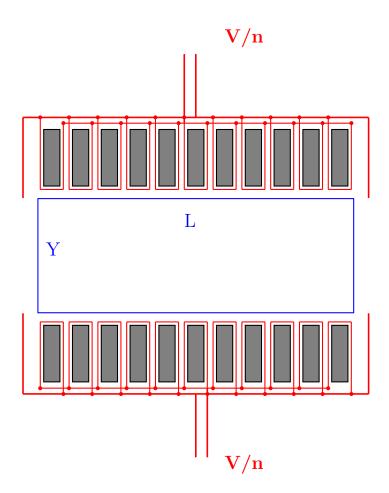
If the load must be made resitive, then the total energy would be 4 times higher (V × I ×  $t_{\text{pulse}}$  vs. V × I ×  $t_{\text{rise}}$  /2).

### Cos Theta Version

- J $\approx \times \pi/4$
- V's not the same for different phi's



End View



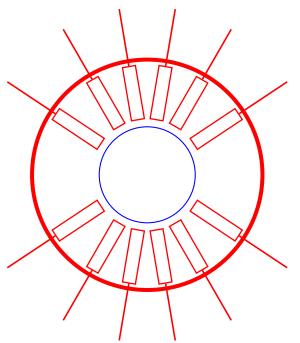
Side View

### Works with no Ferrite

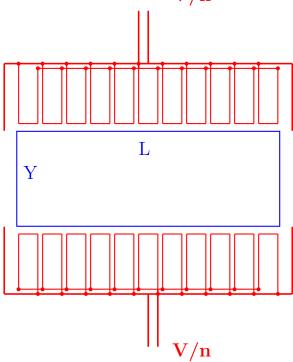
- V = the same
- J $\approx 2 \times$
- I $\approx 2\times$

#### But

- No rise time limit
- ullet Not effected by solenoid fields  $_{
  m V/n}$

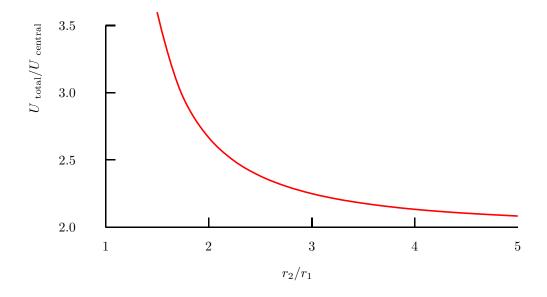


End View



Side View

# Stored Energy vs. rad 2



e.g. 
$$r_2/r_1 = 3$$
:  $U(\text{no Fe}) = 2.25 \times U(\text{Fe})$ 

## "Resonant" Mag Amp Driver

Assume:

- A single compression stage;
- Negligible cable distance from driver to kicker;
- A sudden saturation at  $I_s$  from an initial large inductance  $L_1$  to a small inductance  $L_2$ ;
- A purely inductive kicker magnet with inductance L;
- Start the clock at t=0 with the drive capacitor C charged to  $V_o$  and no current flowing.

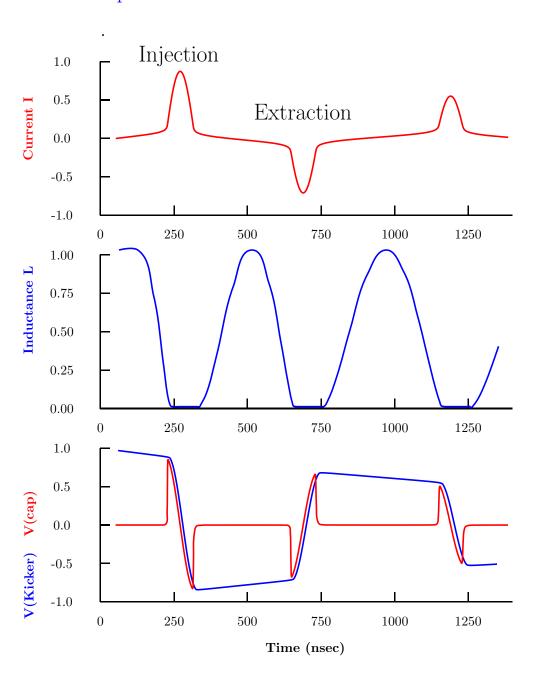
Initially: long time constant  $\tau_L$ 

$$au_L = \sqrt{(L+L_1)C}$$
 $V_{\text{Capacitor}} = V_o \cos\left(\frac{t}{ au_L}\right)$ 
 $V_{\text{Kicker}} = V_1 \left(\frac{L}{L+L_1}\right)$ 
 $I = I_o \sin\left(\frac{t}{ au_L}\right)$ 

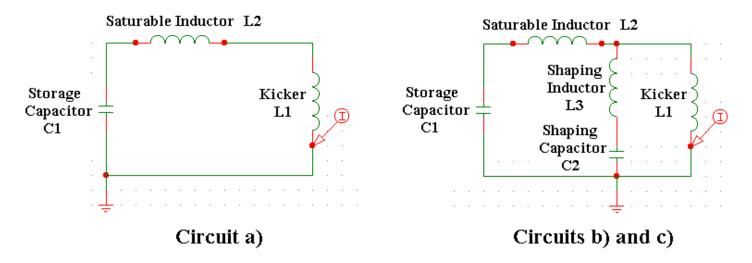
When current = saturation current: Short time constant

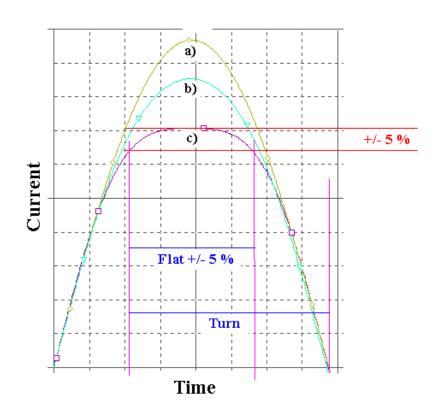
$$\tau_S = \sqrt{(L+L_2)C}$$

# If $\tau_L \approx 100 \times \tau_S$ , Optimistic?



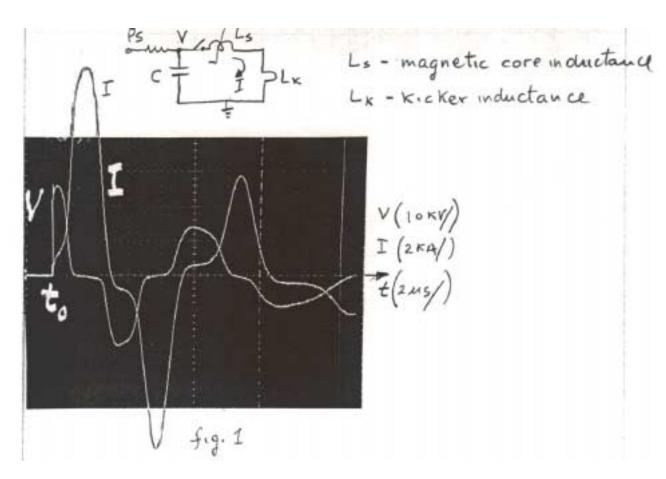
# **Increased Flat Top**





case	C1	L1	L2	C2	L3	V/Va	U/Ua	flat/turn
a)	142	100	7			1.0	1.0	0.46
b)	120	100	7	<b>20</b>	90	0.88	1.29	0.52
<b>c</b> )	100	100	7	<b>40</b>	<b>53</b>	0.73	1.89	0.57

#### Circuit Model



- May be hard for many turns
- Not simulated with initial stages

#### Conclusion

This study suggests that even the first cooling ring injection and extraction problems may be soluble.

But much work remains.

The proposed kicker could also have application in FFAG acceleration of large emittance beams for neutrino factories.