

# Introduction

We'll see that:

- Laser cooling, by Compton scattering, works very well in theory
  - Power needed scales as  $1/E^2$  of muon beam
- ⇒ Might be feasible for  $\geq 1.5$  TeV beam  
 $\geq \times 1000$  increase in luminosity per muon relative to parameters in status report

Scenario:

~~Go directly to~~ 3 TeV c.o.m.

*COLLIDER IN STATUS REPORT*

Still need front-end cooling, but looser requirements:

- Lower repetition rate
- Less cooling— just good enough to get into accelerator
- Lower muon acceptance

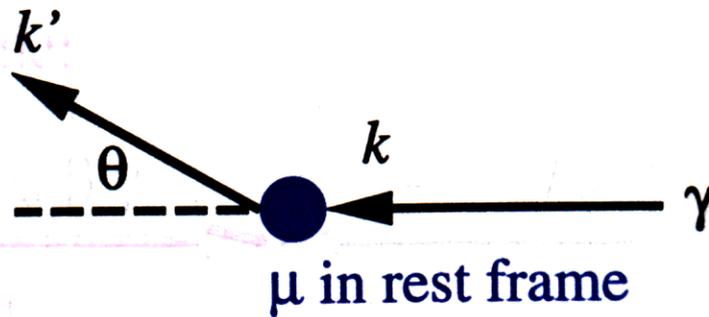
Other benefits

- Reduced radiation problem
- Reduced detector backgrounds
- Possibility for ultra-high luminosity

New set of technical miracles needed, but:

- Net improvement in feasibility?
- Net reduction in power consumption?

# Compton Scattering



$$\frac{d\sigma}{d\Omega} = \left( \frac{\alpha^2}{2m^2} \right) \left( \frac{k'}{k} \right)^2 \left( \frac{k'}{k} + \frac{k}{k'} - \sin^2\theta \right)$$
$$k' = \frac{k}{1 + (k/m)(1 - \cos\theta)}$$

For  $k \ll m$ :

$$k' \approx k$$

$$\frac{d\sigma}{d\Omega} = 1 + \cos^2\theta$$

$$\sigma = \frac{8\pi}{3} r_\mu^2$$

$r_\mu^2$  is the classical muon radius,  $\frac{\alpha^2}{m^2}$

$\sigma$  reduced by  $4 \times 10^4$  relative to  $e$

For 0.1 eV  $\gamma$ : 1 collision  $\Leftrightarrow 10 \text{ J} / (\mu\text{m})^2$

# Laser Cooling of 1.5 TeV Muon Beams

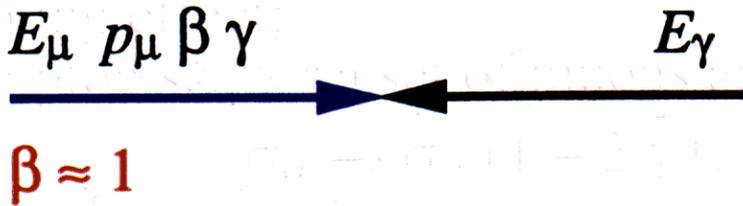
Fritz DeJongh

- Introduction
- Compton scattering
- Transverse cooling effect
- Longitudinal cooling effect
- Configuration with storage ring
- Conclusions

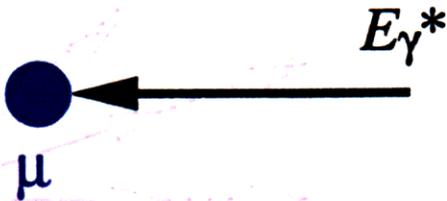
*Application to  $e^+e^-$  beams  
for  $\gamma\gamma$  collider*

*Palmer  
Telnov  
Ohsugi*

Lab frame:



$\mu$  frame:



$$E_\gamma^* = \gamma E_\gamma (1 + \beta) \approx 2 \gamma E_\gamma$$

- On average:
- In  $\mu$  frame:

$\gamma$  transfers energy  $E_\gamma^*$  to  $\mu$

Average  $E_\mu$  loss in lab frame:

$$E_\mu \rightarrow E_\mu - 2 \gamma^2 E_\gamma$$

Average transverse kick:

$$p_{T\mu} \sim \gamma E_\gamma$$

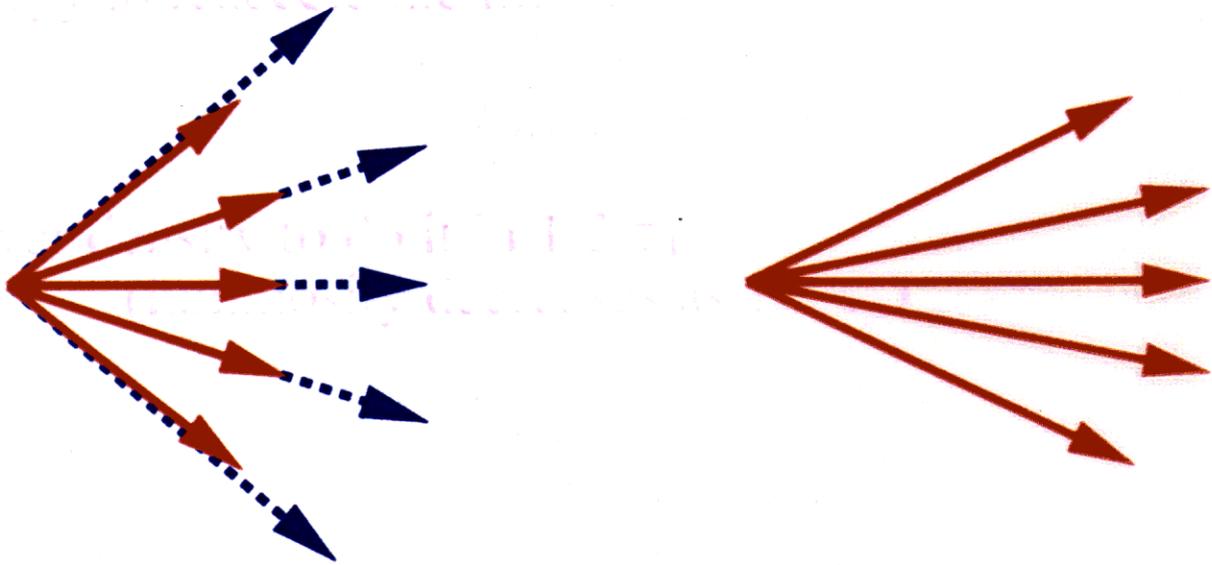
$\Rightarrow$  For large  $\gamma$ ,  
 $\mu$  slows down without changing direction.

# Transverse Cooling Effect

Cooling effect:

Compton scatter a set of muons and reaccelerate:

$$\sigma_{\alpha} \rightarrow \sigma_{\alpha} (1 - 2 \gamma E_{\gamma} / m)$$



Heating effect:  
From transverse kick

$$\sigma_{\alpha}^2 \rightarrow \sigma_{\alpha}^2 + 2 \frac{15}{16} \left( \frac{E_{\gamma}}{m} \right)^2$$

Poisson stat

Equilibrium:

$$\sigma_{\alpha} = \left( \frac{15}{32} \frac{E_{\gamma}}{E_{\mu}} \right)^{1/2}$$

## To decrease $\sigma_\alpha$ by $1/e$

Number of scatterings

$$n = \frac{m}{2\gamma E_\gamma}$$

Energy to reaccelerate muon

$$E_{reacc} = E_\mu$$

Power density to do it in  $1/2 \tau_\mu$   
(luminosity decreases as  $1/2 \tau_\mu$ )

$$p = \frac{3m}{8\pi \gamma^2 \tau_\mu r_\mu^2}$$

Decreases as the square of  $E_\mu$  increases!

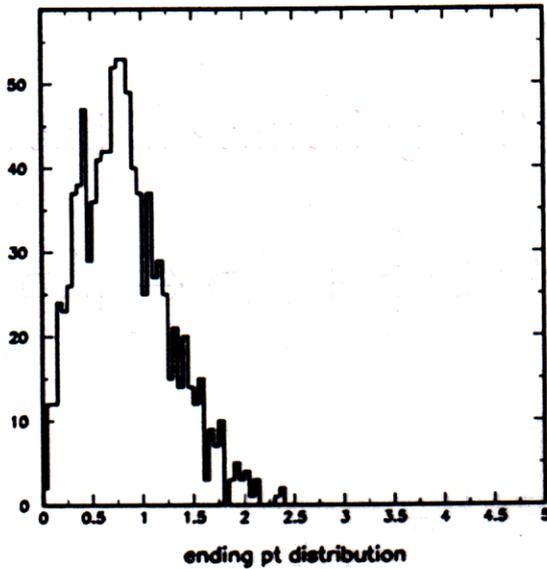
For  $E_\mu = 1.5 \text{ TeV}$ :

$$p = 2.4 \times 10^7 \text{ W} / \mu\text{m}^2$$

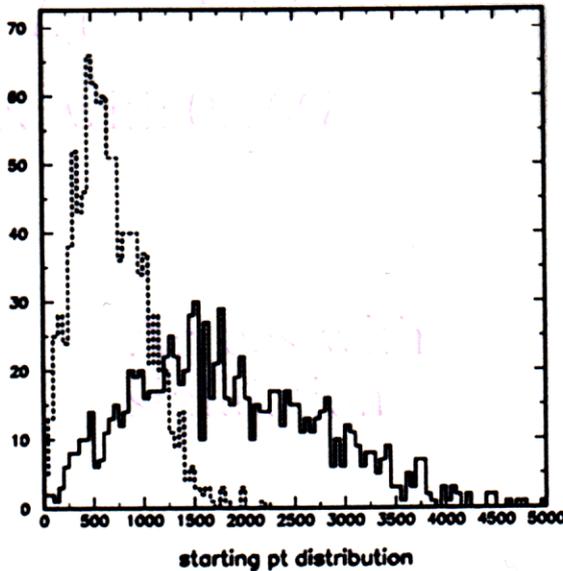
# Simple simulation – quantitative agreement

$$E_\gamma = 1 \text{ eV}, E_\mu = 1.5 \text{ TeV} \Rightarrow \sigma_\alpha = 5.6 \times 10^{-7} \text{ rad}$$

compared with  $10^{-3} \text{ rad}$



Equilibrium  $p_T$



$p_T$  reduction

Dashed: ending  $p_T$  distribution

# Storage Ring Longitudinal Cooling Effect

Energy loss  $\propto E_\mu^2$

$$\sigma_{E_\mu} \rightarrow \sigma_{E_\mu} - \frac{\sigma_{E\mu}}{E_\mu} 4\gamma^2 E_\gamma$$

RMS spread in scattering + Poisson term

$$\sigma_{E_\mu}^2 \rightarrow \sigma_{E_\mu}^2 + \frac{2}{5} (2\gamma^2 E_\gamma)^2 + (2\gamma^2 E_\gamma)^2$$

Equilibrium

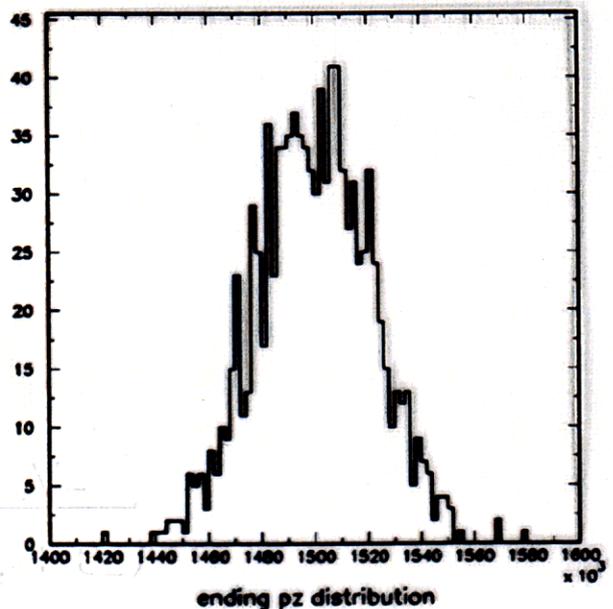
$$\frac{\sigma_{E\mu}}{E_\mu} = \left( \frac{7\gamma E_\gamma}{5m} \right)^{1/2}$$

For  $E_\gamma = 1 \text{ eV}$ ,  $E_\mu = 1.5 \text{ TeV}$

$$\frac{\sigma_{E\mu}}{E_\mu} = 1.4\%$$

compared with 0.16%

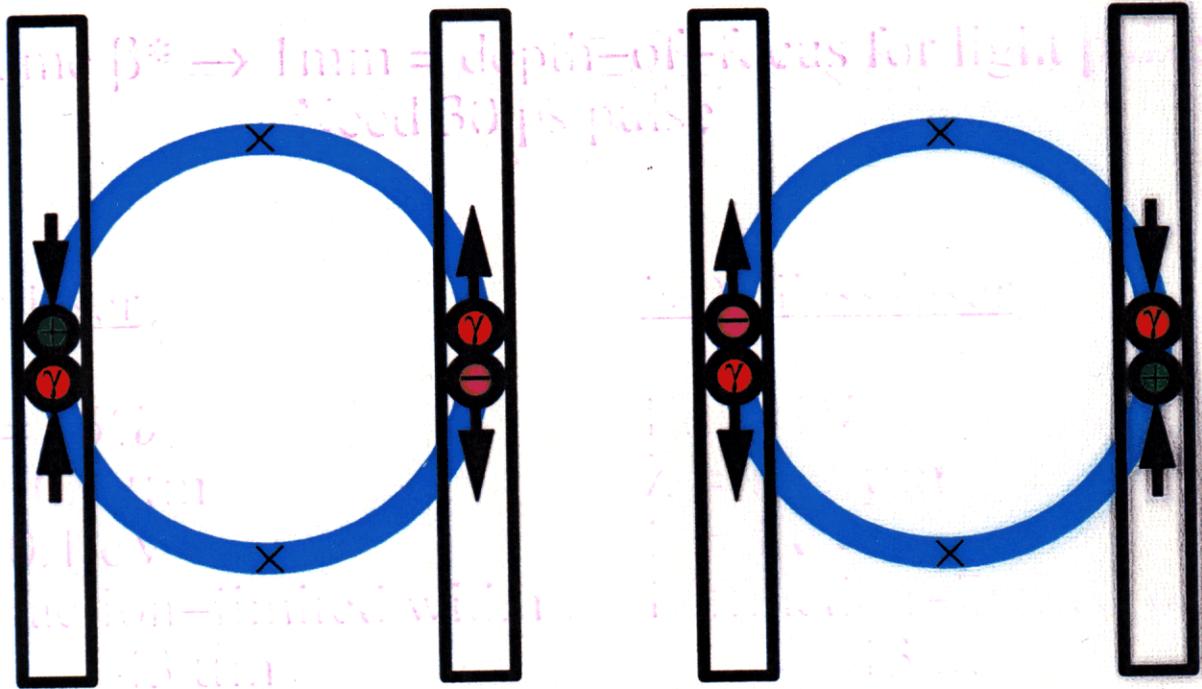
Agrees with simulation



# Use with 1.5 TeV/beam Storage Ring

- 1.5 TeV
- 750 turns

⇒ Assume 4 bunches → 1 bunch



Parameters from status report:

- circumference = 6 km
- $\beta^* = 3 \text{ mm}$
- $\sigma_z = 3 \text{ mm}$
- $\sigma_r = 3.2 \text{ } \mu\text{m}$
- $\sigma_\theta = 1.1 \text{ mrad}$
- $L = 7 \times 10^{-34} / \text{cm}^2 \text{ sec}$

$$L \propto \frac{N_+ N_-}{\beta^* (\epsilon_+ + \epsilon_-)}$$

To cool in  $1 \tau_L$ :

$(1/e)$

- 15 msec
- 750 turns

$\Rightarrow Q = 1000$ , reasonable

Assume  $\beta^* \rightarrow 1 \text{ mm} = \text{depth-of-focus for light pulse}$

Need 30 ps pulse

$p \propto \beta^*$

OK IF  $\sigma_E$  INCREASES

CO<sub>2</sub> laser

ND:Glass laser

Eff = 25%

Eff = 1%

$\lambda = 10.6 \mu\text{m}$

$\lambda = 1.05 \mu\text{m}$

E = 0.1 eV

E = 1 eV

Diffraction-limited width

Diffraction-limited width

43  $\mu\text{m}$

13  $\mu\text{m}$

F-stop = 7

F-stop = 22

Lum factor: 5600

Lum factor: 1770

Energy spread: 0.4%

Energy spread: 1.4%

Energy in pulse:

Energy in pulse:

0.5 MJ

0.05 MJ

x 8.6 FOR eq.

x 7.5 FOR eq.

# Conclusions

>  $\times 1000$  in luminosity  
~10 MW of power

$\Rightarrow$  Can in principle be done with reasonable power consumption, and result in looser requirements on other components.

CO<sub>2</sub> laser seems preferable

- Better cooling
- Lower energy spread

Can a laser be made with

- short enough pulses?
- high enough power

Can good enough optics be made?

Can we get the light to the IP?

Can we make muon bunches shorter?